

Anomalous Hall/Nernst Effects in Magnetic Semiconductors and Magnetic Insulators

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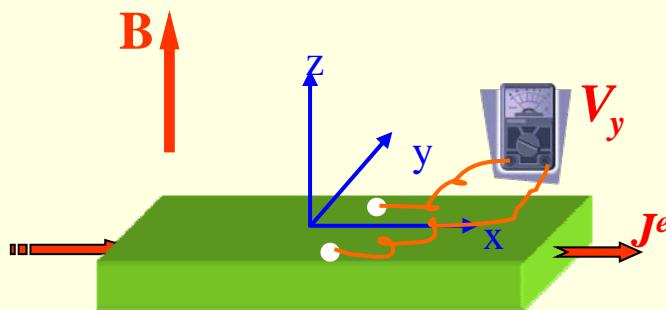
Outline

- Introduction
 - Anomalous Hall effect (AHE)
 - Anomalous Nernst effect (ANE)
- GaMnAs: Dilute Magnetic Semiconductor (DMS)
 - AHE/ANE in absence of B-field
 - Validity of Mott relation with $n=2$
 - Hole-mediated ferromagnetism, probably via impurity band
- Fe_3O_4 : Ferrimagnetic Insulator (FMI)
 - Robust AHE power-law scaling with $n=0.3$
 - Preliminary ANE data: absence of ANE
- Summary

Anomalous Hall Effect (AHE)



In ferromagnets, ρ_{xy} contains two parts:



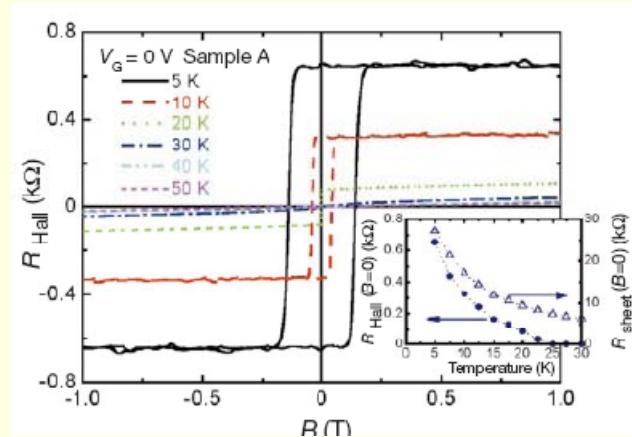
Normal or ordinary Hall effect (OHE)

$$\rho_{xy} = R_0 B + R_s M$$

ρ_{AH} : anomalous or
extraordinary Hall effect (AHE)

- ρ_{xy} exists even if $B=0$
- AHE is more than an order greater than OHE
- $\rho_{xy} \sim \rho^{AH}$

AHE is not caused by magnetic field, but by spin-orbit coupling (SOC)



AHE from 5 nm-thick (In, Mn)As layer

Physical Origin of AHE

Spin-orbit effect: extrinsic (scattering) or intrinsic (band structure)

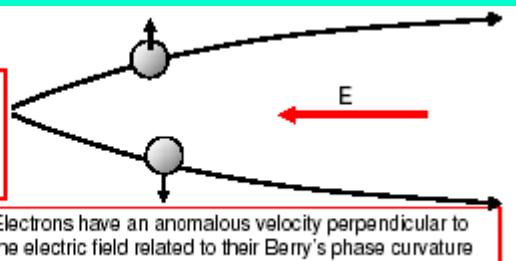
❖ Karplus & Luttinger (intrinsic: inter-band effect)

Intrinsic deflection

Interband coherence induced by an external electric field gives rise to a velocity contribution perpendicular to the field direction. These currents do not sum to zero in ferromagnets.

$$\frac{d\langle \vec{r} \rangle}{dt} = \frac{\partial E}{\hbar \partial \vec{k}} + \frac{e}{\hbar} \vec{E} \times \vec{b}_v$$

Electrons have an anomalous velocity perpendicular to the electric field related to their Berry's phase curvature

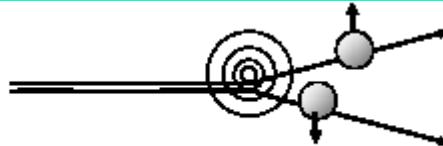


See excellent review articles: N.A. Sinitsyn, J. Phys.: Condens. Matter (2008); N. Nagaosa et al., Rev. Mod. Phys. (2010).

❖ Smit (extrinsic: skew scattering)

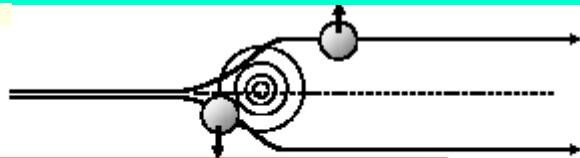
Skew scattering

Asymmetric scattering due to the effective spin-orbit coupling of the electron or the impurity.



❖ Berger (extrinsic: side-jump)

Side jump



The electron velocity is deflected in opposite directions by the opposite electric fields experienced upon approaching and leaving an impurity. The time-integrated velocity deflection is the side jump.

❖ Niu & MacDonald (intrinsic: Berry's phase)

Power-Law

Power-law:

$$R_s = \lambda \rho_{xx}^n$$

❖ Exponent n=2:

$$\sigma_{xy} \sim \rho_{xy} / \rho_{xx}^2 \quad (\rho_{xx} \gg \rho_{xy})$$

→ independent of $1/\tau$!

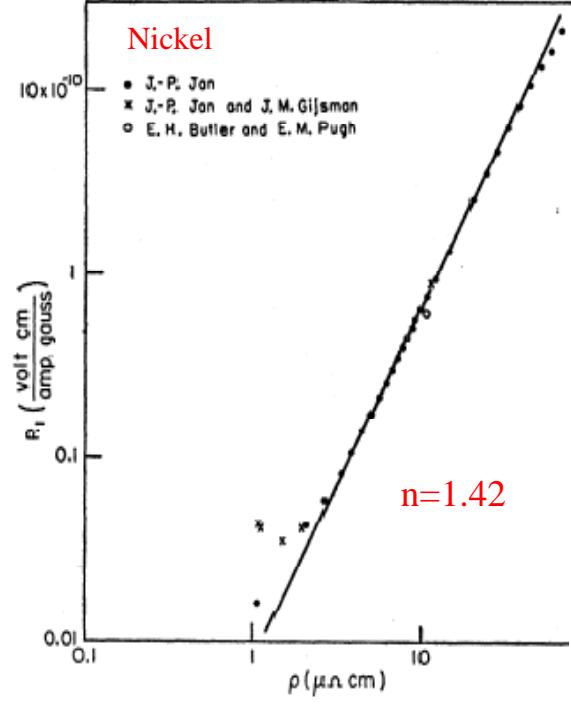
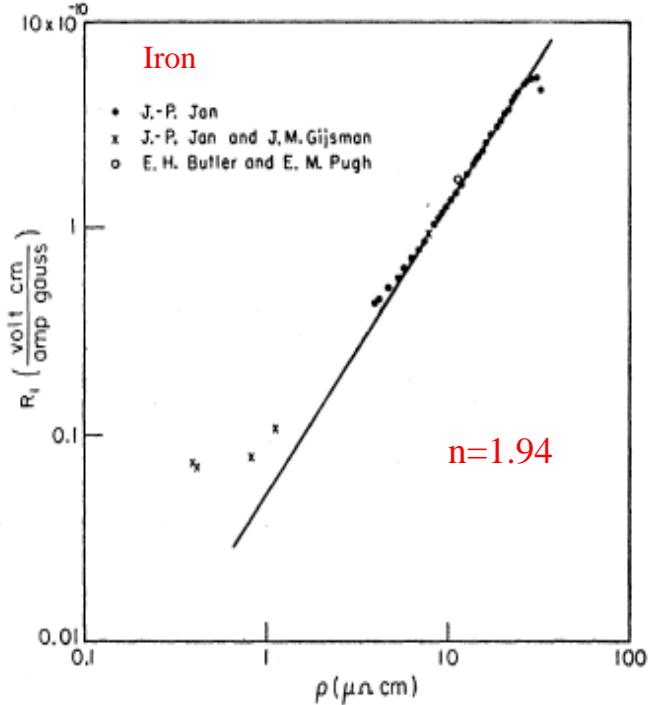
→ Special Hall current: $J_H = \sigma_{xy} E_x$

❖ Exponent n=1:

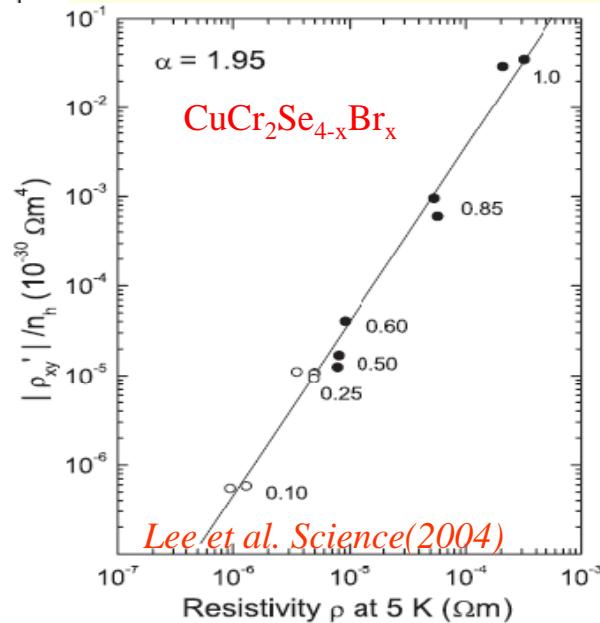
$$\sigma_{xy} \sim \sigma_{xx}$$

skew scattering (extrinsic)

Experimental Data



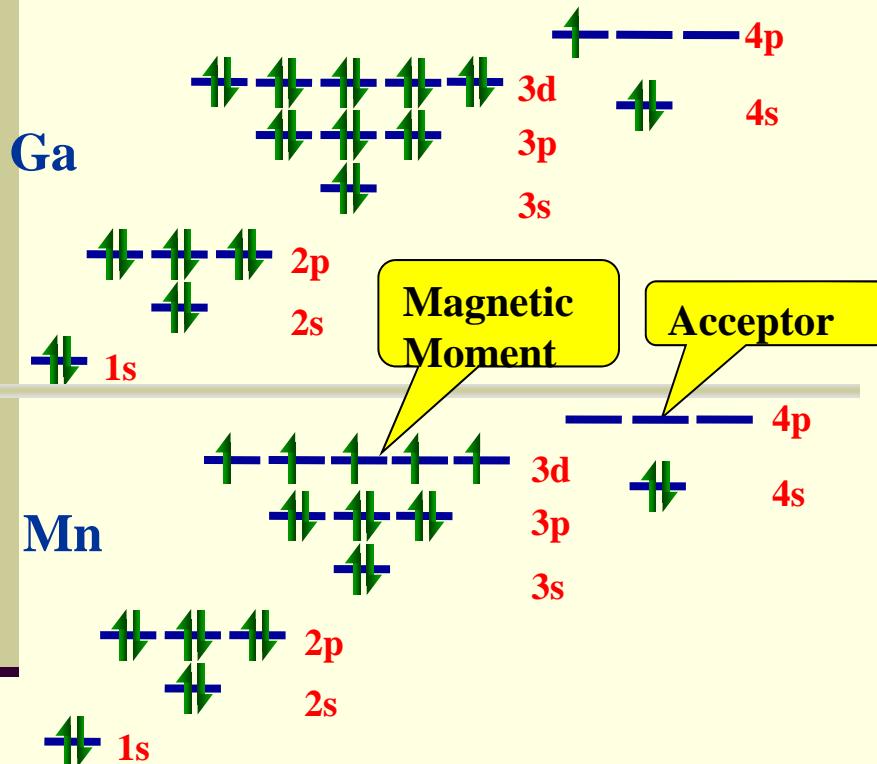
Pure metals are not ideal



- Scattering rate can be tuned by impurity, temperature or magnetic field
- Semiconductors or alloys are preferred

Dilute Magnetic Semiconductors (DMS)

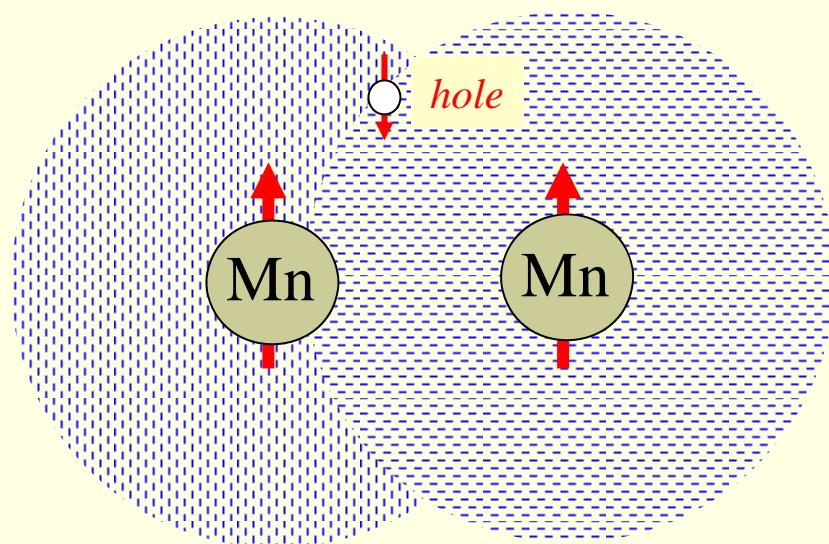
$\text{Ga}_{1-x}\text{Mn}_x\text{As}$: most studied DMS



Mn substitutes Ga in GaAs: introducing spin and charge carriers!

■ Carrier-mediated interaction

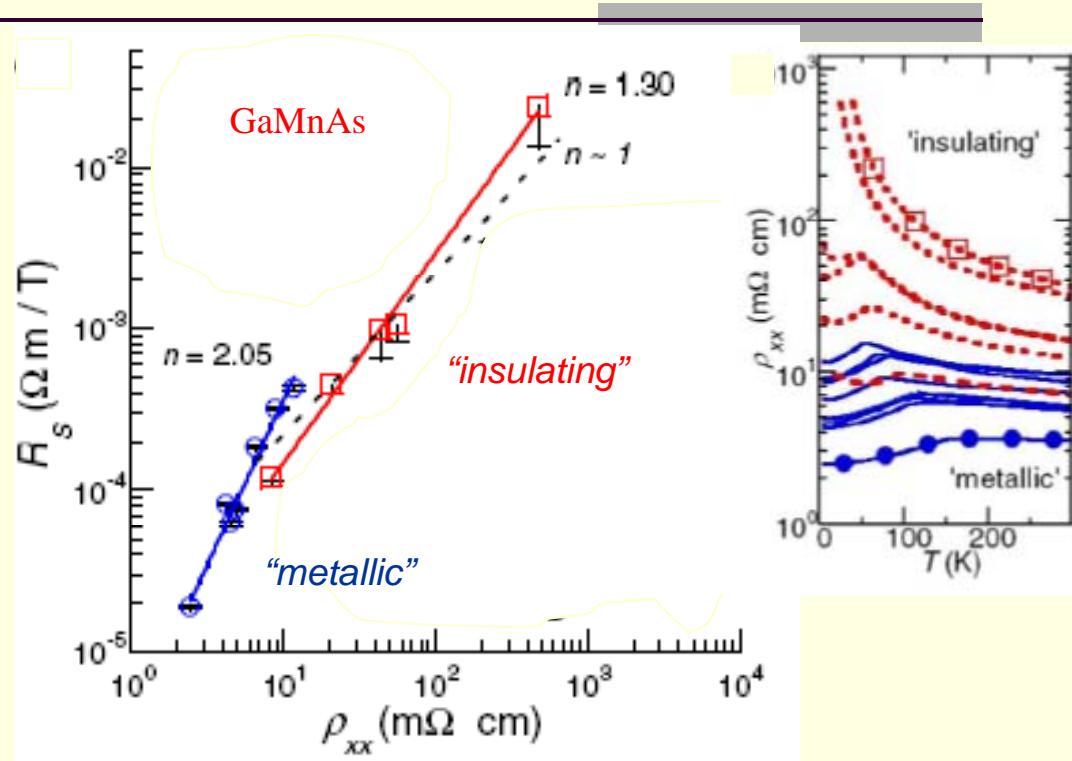
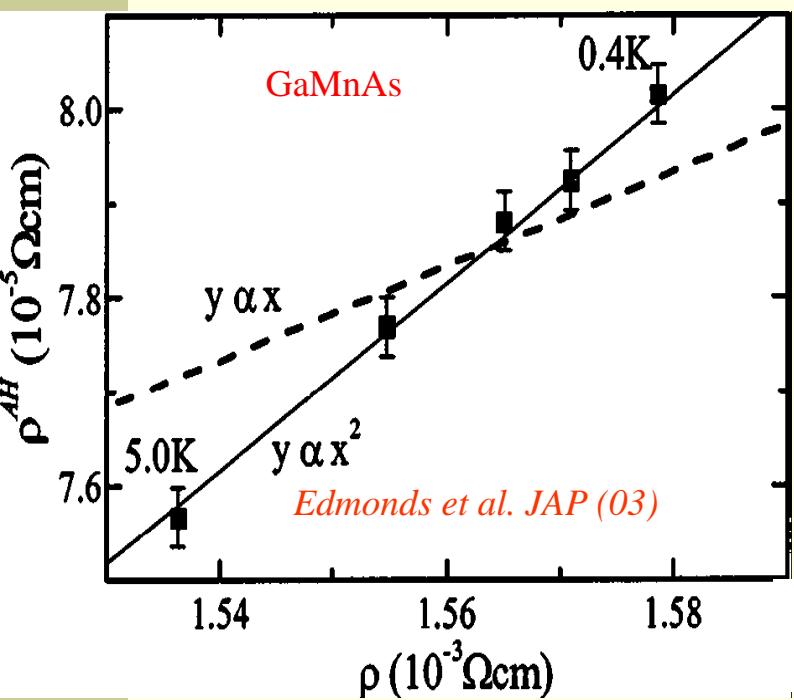
Zener model: $T_c \sim x^*p^{1/3}$



Highest $T_c \sim 150$ K!

- Strong SOC for holes in GaAs
- Strong impurity scattering

AHE in DMS



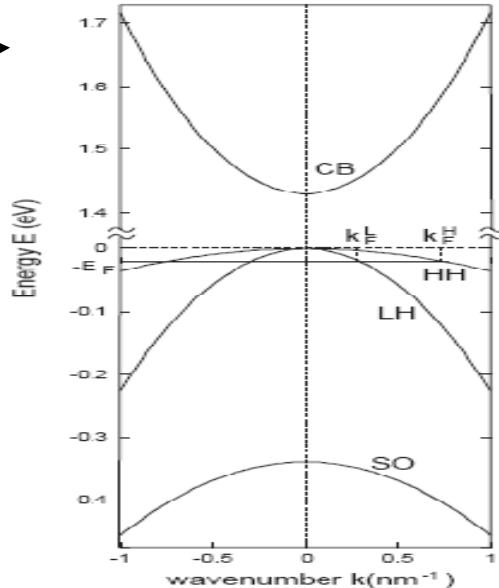
Results support intrinsic mechanism (but the resistivity range is too narrow)

$n \sim 2$ and AHE's carrier density dependence in "metallic" regime \rightarrow intrinsic origin

DMS: Intrinsic AHE

GaAs band structure →

- Spin-orbit coupling
- External electric field E



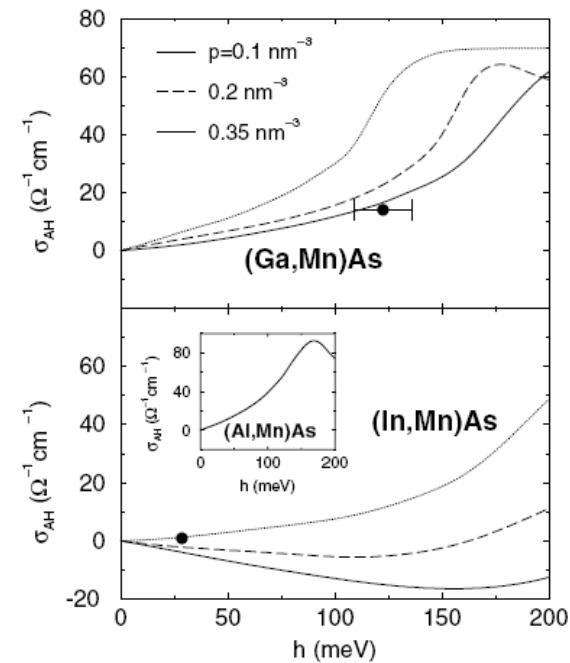
Electron wavepackets acquire additional velocity:

$$\dot{\vec{x}}_c = \frac{\partial \epsilon}{\hbar \partial \vec{k}} + \frac{(e/\hbar) \vec{E} \times \vec{\Omega}}{}$$

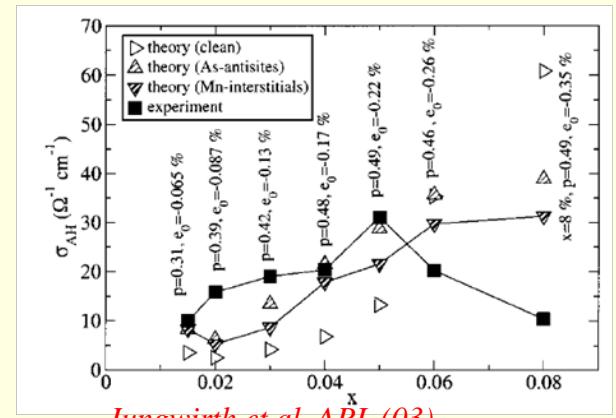
Anomalous velocity

Under broken time reversal symmetry, this Berry phase effect alone gives rise to AHE comparable with experimental values

→ Intrinsic origin of AHE in DMS

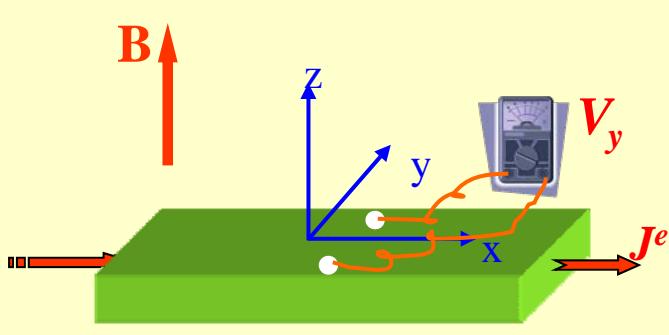


Jungwirth, Niu and McDonald, PRL (02)



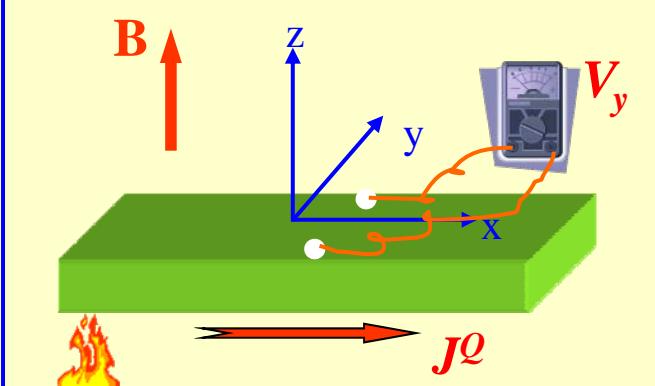
Jungwirth et al. APL (03)

Nernst Effect



Hall effect

$$\sigma_H = \frac{E_y}{J_x}$$



Nernst effect

$$S_N = \frac{E_y}{(\nabla T)_x} = - \frac{\Delta V_y}{\Delta T_x} \cdot \frac{L}{W}$$

- Normal Hall Effect

$$\sigma_H \propto B$$

- Anomalous Hall Effect (AHE)

$$\sigma_{AH} \propto M$$

- Normal Nernst Effect

$$S_N \propto B$$

- Anomalous Nernst Effect (ANE)

$$S_{AN} \propto M$$

Other Transport Effects

Stimulus: ΔV_x or ΔT_x

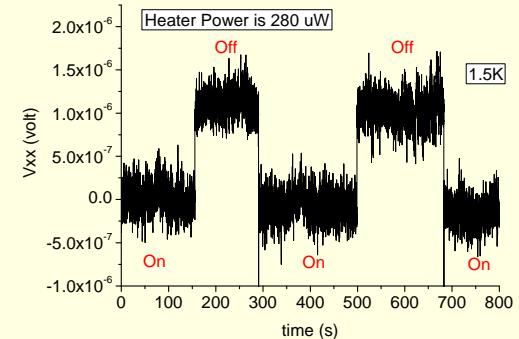
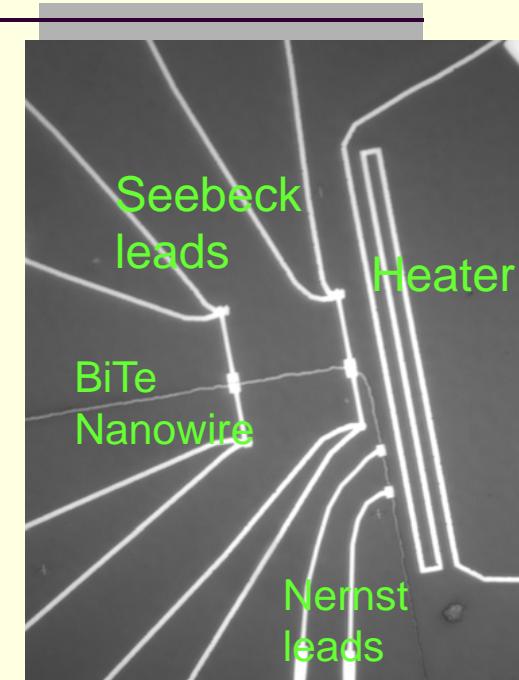
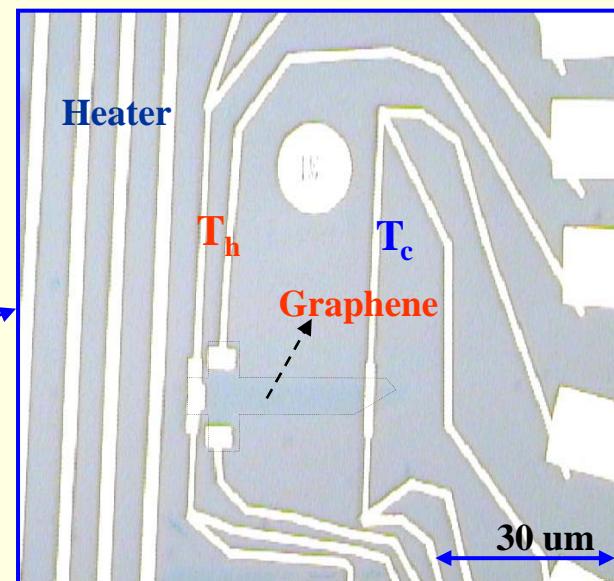
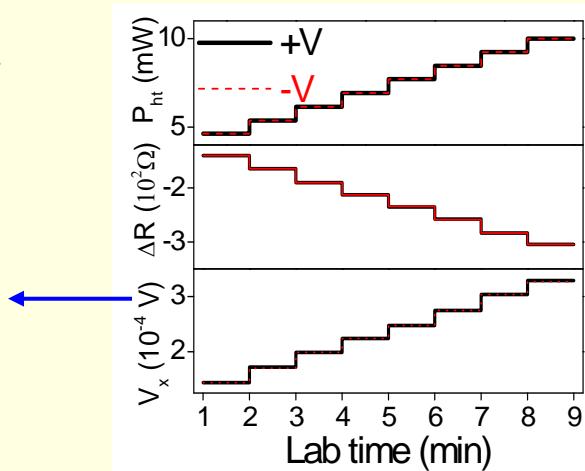
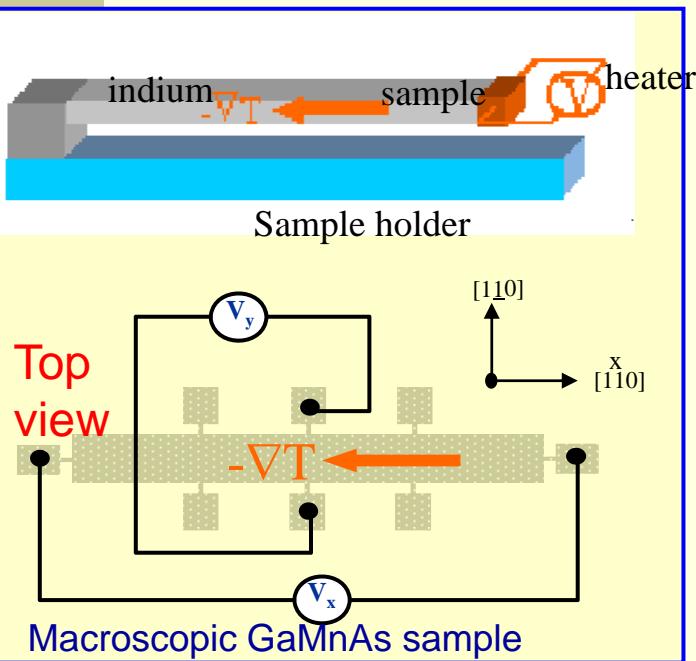
Response: $\rightarrow I_x, I_x^Q$ and I_x, I_x^Q if there is B or M.

Open-circuit condition: measuring ΔV instead of I; ΔT instead of I^Q

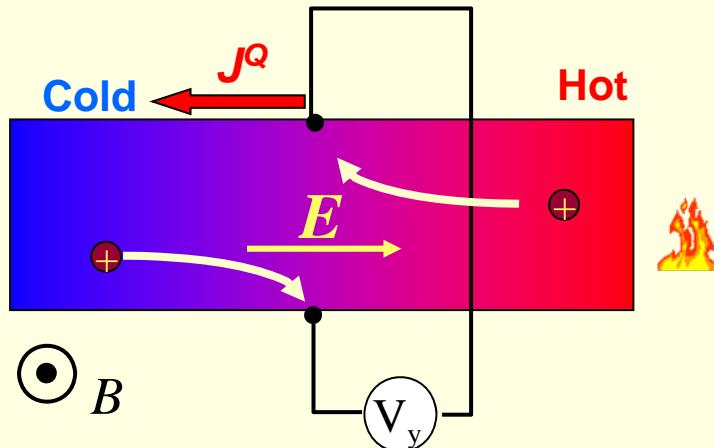
Stimulus \ Response	I_x	I_y	I_x^Q or ΔT_x	I_y^Q or ΔT_y
ΔV_x	Conductivity	Hall effect	Peltier effect	Ettingshausen effect
ΔT_x	Seebeck effect	Nernst effect	Thermal conductivity	Righi-Leduc effect

Various coefficients are connected by Onsager relations and other relations (e.g. Wiedemann-Franz law, Mott relation, etc.)

Thermoelectric Measurements



Nernst Effect



Open circuit voltage (J_x and $J_y=0$)

$$\vec{J} = \sigma \vec{E} + \alpha(-\nabla T)$$

$$S_{yx} = \frac{1}{\sigma_{xx}} (\alpha_{yx} - \sigma_{yx} S_{xx})$$

diffusion *drift*

Nernst current *Hall current*

- Nernst effect is a net effect resulting from both Hall current (drift) and Nernst current (diffusion). In metals, it is a very small effect.

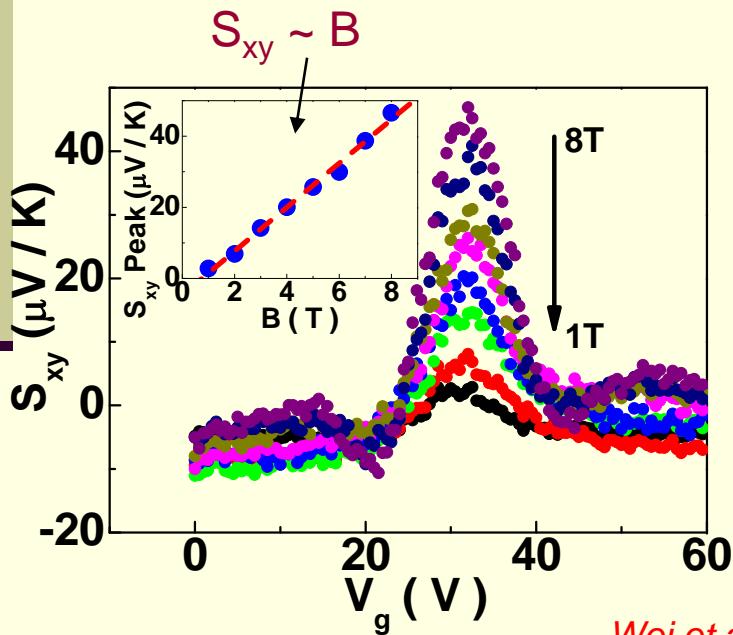
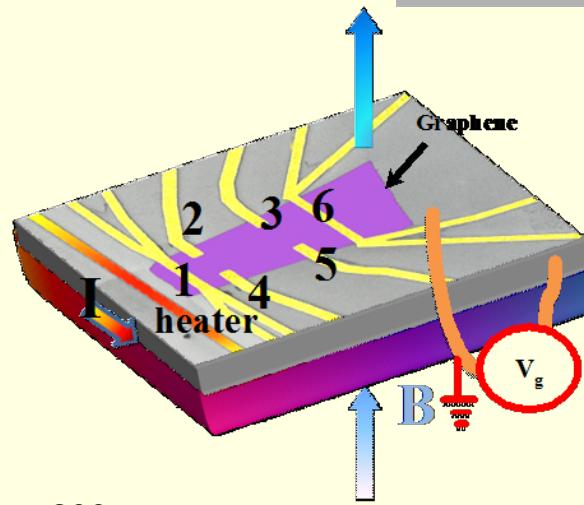
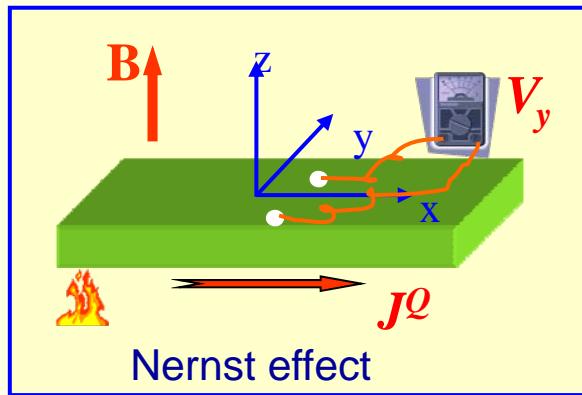
Mott relations

$$\alpha_{xx} = \frac{\pi^2 k_B^2 T}{3e} \left(\frac{\partial \sigma_{xx}}{\partial \epsilon} \right)_{\epsilon_F}$$

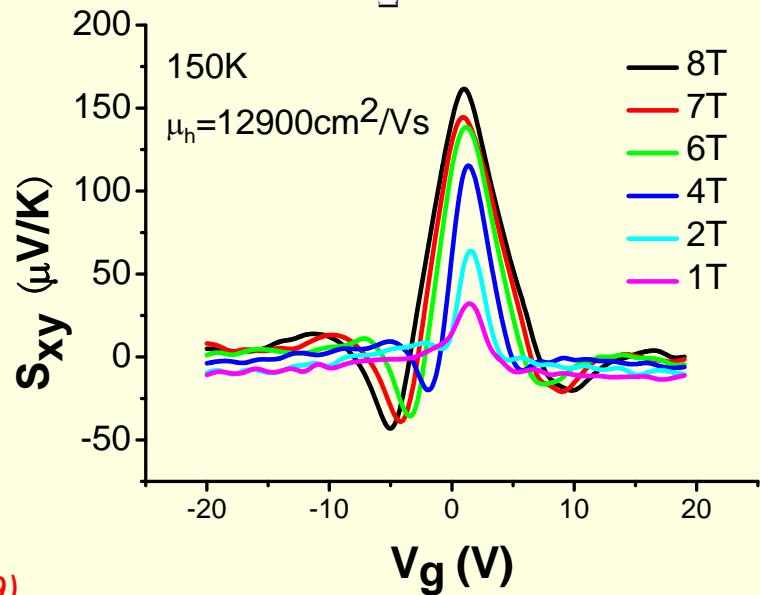
$$\alpha_{yx} = \frac{\pi^2 k_B^2 T}{3e} \left(\frac{\partial \sigma_{yx}}{\partial \epsilon} \right)_{\epsilon_F}$$

$$S_{yx} = \frac{\pi^2 k_B^2 T}{3e} \left[\arctan \left(\frac{\sigma_{yx}}{\sigma_{xx}} \right) \right]' = \frac{\pi^2 k_B^2 T}{3e} \left(\frac{\partial \Theta_H}{\partial \varepsilon} \right)_{\varepsilon_F}$$

Normal Nernst Effect in Graphene



Wei et al. PRL (09)



Why Anomalous Nernst Effect?

- Does ANE exist if AHE is caused by intrinsic effect?
 - Only statistical force (no E-field)

$$\dot{x}_c = \frac{\partial \epsilon}{\hbar \partial k} + (e/\hbar) \vec{E} \times \vec{\Omega}.$$

- If \dot{x}_c does not depend on Fermi energy, there will be no ANE

$$\rho_{AH} = \lambda \rho_{xx}^2 M_z \rightarrow \sigma_{AH} = \lambda M_z \quad \text{Finite } S_{AN} \rightarrow \lambda = \lambda(\epsilon_F)!$$

- How are AHE and ANE related?
 - Validity of Mott relation (D. Xiao PRL 07)
 - Does ANE help us understand physical origin of AHE
- What can we learn about ferromagnets?
 - Spin-orbit coupling, magnetic ordering

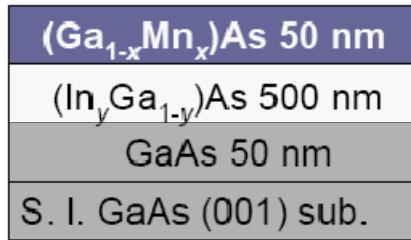
Subtleties with DMS

- Films with in-plane anisotropy often require high magnetic fields to obtain finite AHE (i.e. M_z). To completely saturate M_z , it requires magnetic fields in excess of 10 T!
- Normal Hall/Nernst signals become large at high fields.
- High magnetic fields cause significant magneto-resistance (i.e. change in ρ_{xx}), even when AHE saturates.
- In DMS, all spins contribute to magnetization, but only those in hole-rich regions contributes to AHE; therefore, it is difficult to separate these two for $\rho_{AH} = \lambda \rho_{xx}^2 M_z$.

Solutions:

- DMS films with perpendicular anisotropy (no need to have B-field)
- ANE and AHE measured simultaneously from the same area (no need to measure M)

GaMnAs with Perpendicular Anisotropy

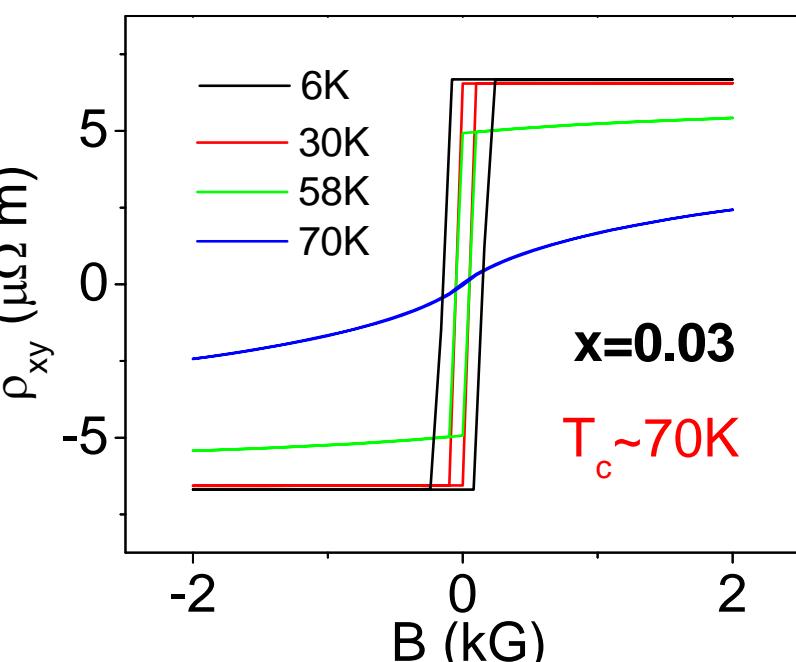


InGaAs buffer layer → tensile strain
→perpendicular anisotropy

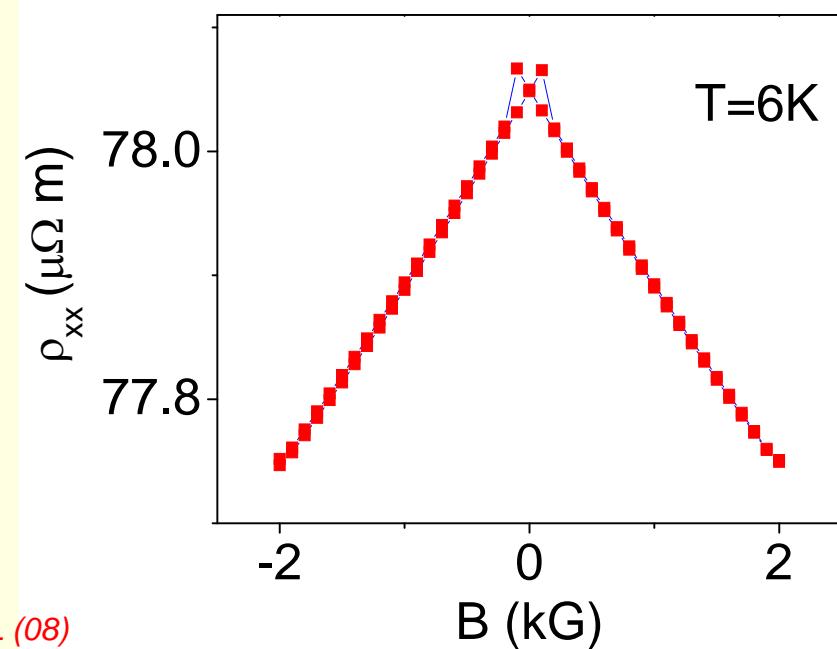
“B=0” is a special field

A set of films with different Mn, so different T_c ’s.

AHE (OHE is negligible)



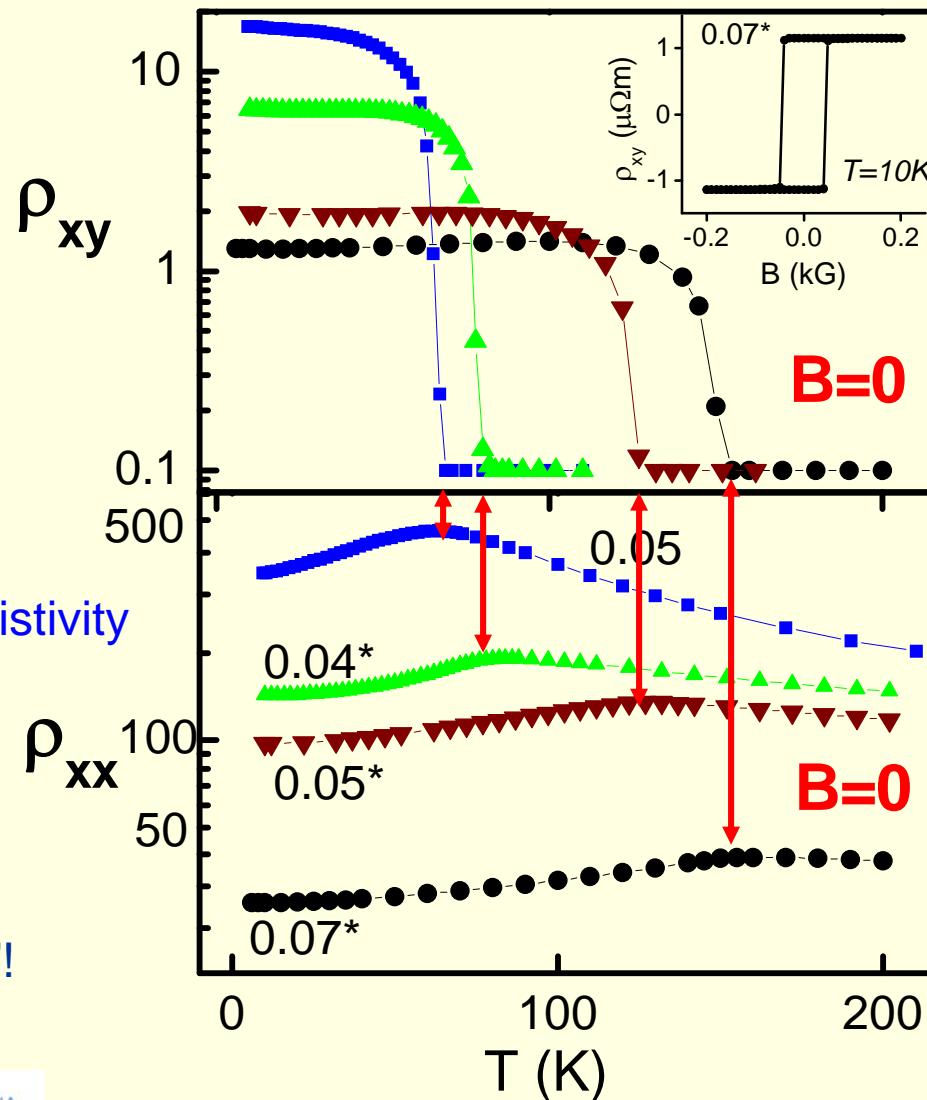
Longitudinal resistivity



Pu et al. PRL (08)

AHE in GaMnAs

Hall resistivity



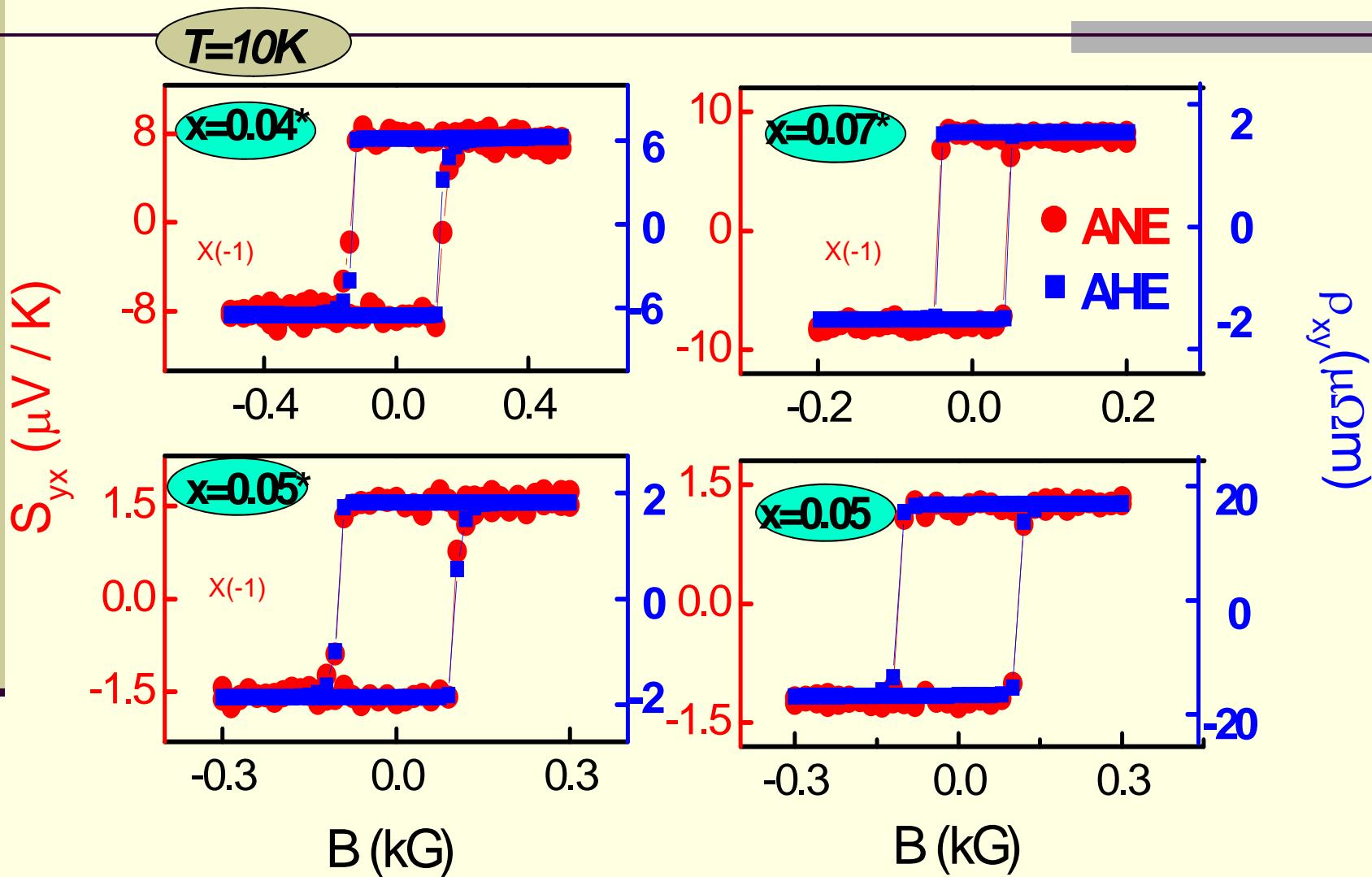
AHE remains positive below T_c

$x=0.04^*$,
 0.05^* , and
 0.07^* are annealed but
 $x=0.05$ is not.

All "metallic"!

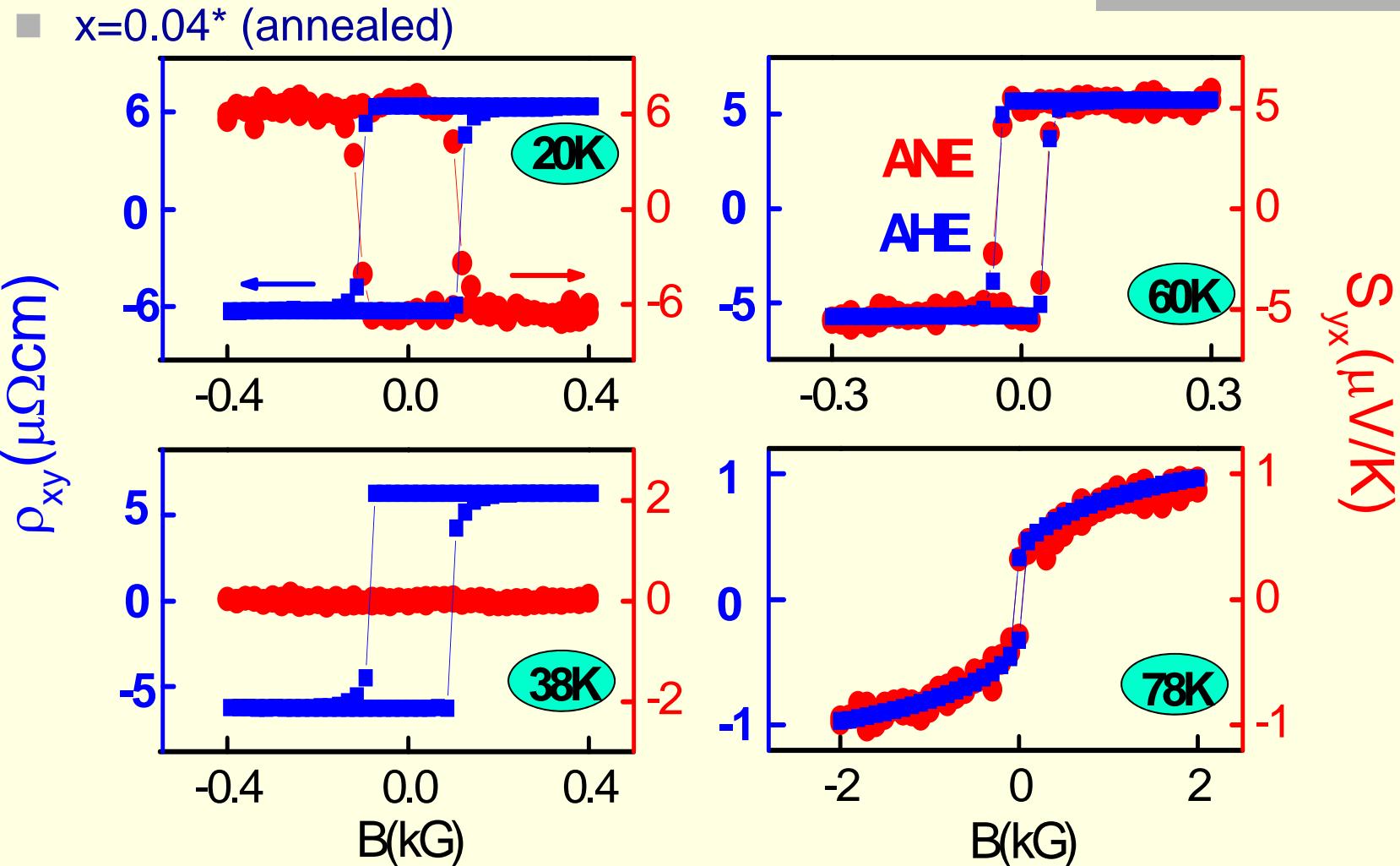
Usual power-law analysis is highly unreliable because of uncertainty in determining M

AHE & ANE in Different Samples



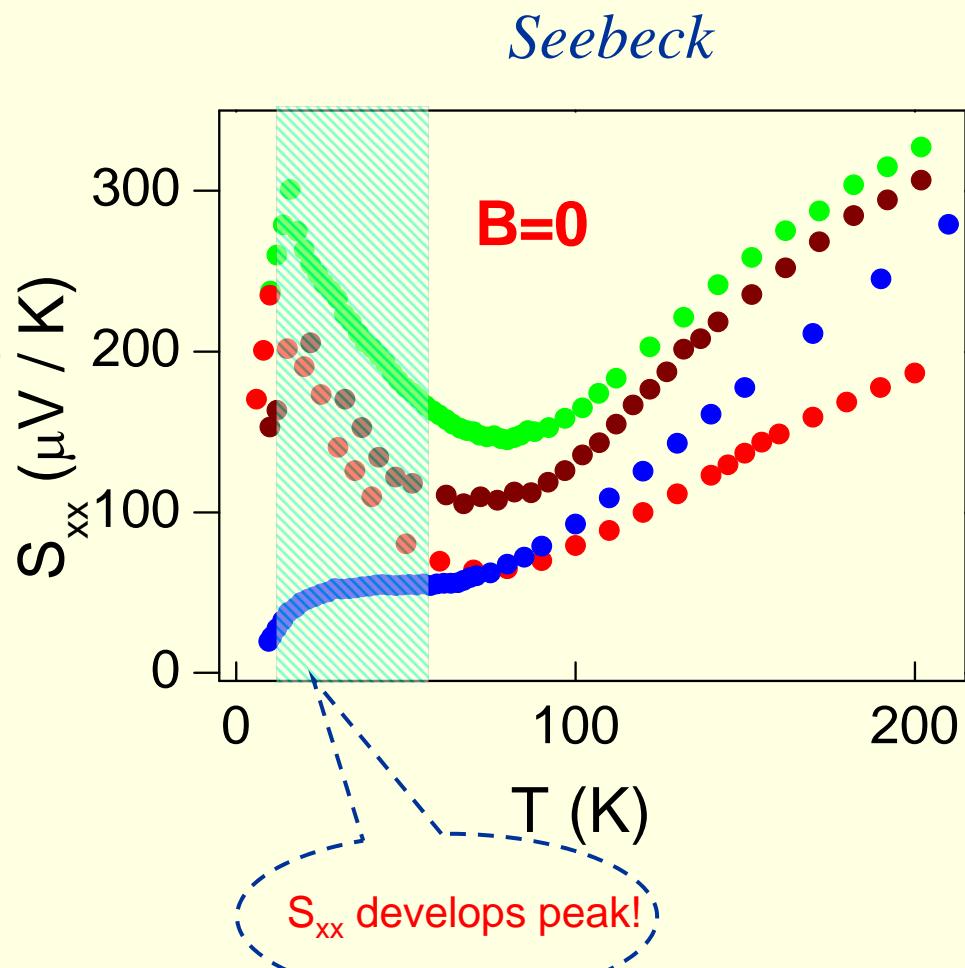
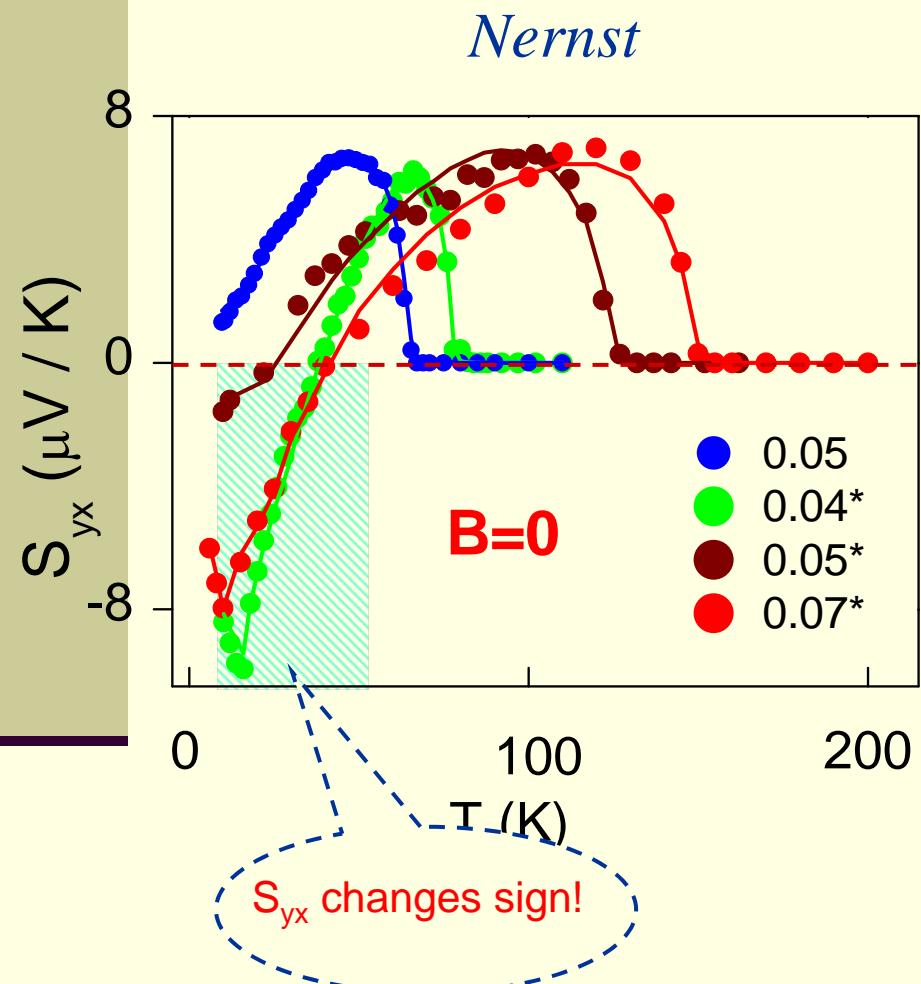
No ordinary Nernst effect is visible; S_{yx} goes with ρ_{xy}

AHE & ANE at Different Temperatures



- Except for sign, AHE and ANE scale with each other → share the same physical origin!

Zero-Field S_{xx} and S_{yx}



Mott Relation

$$\rho_{xy} = \lambda M \rho_{xx}^n$$

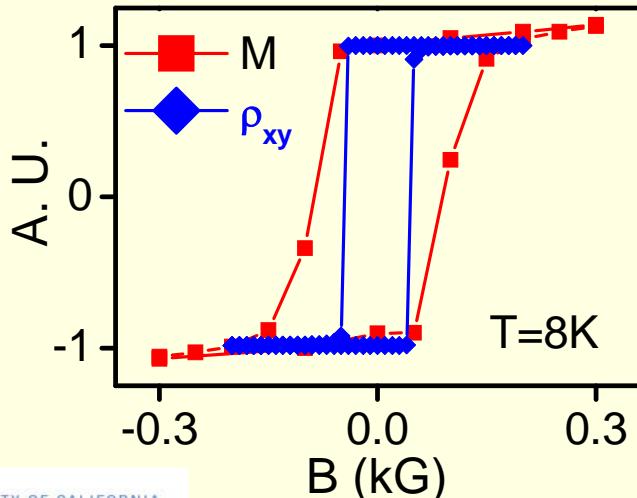
Mott relation \rightarrow

$$S_{yx} \propto \left(\frac{\rho_{xy}}{\rho_{xx}} \right)^n = \lambda M \rho_{xx}^{n-1} \left(T \frac{\pi^2 k_B^2}{3e} \frac{\lambda'}{\lambda} - (n-1) S_{xx} \right)$$

We have

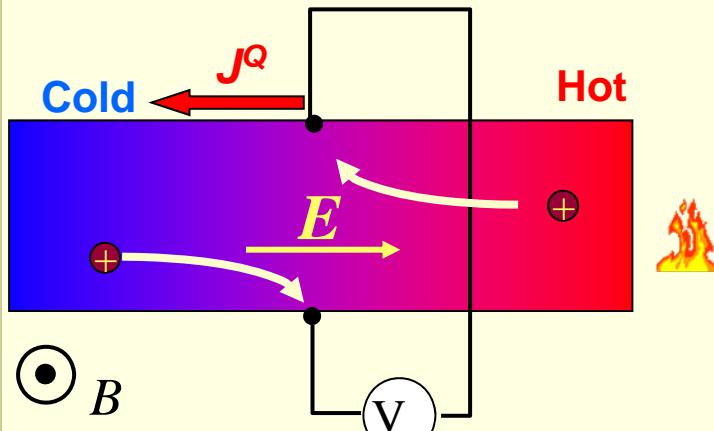
$$S_{yx} = \frac{\rho_{xy}}{\rho_{xx}} \left(T \frac{\pi^2 k_B^2}{3e} \frac{\lambda'}{\lambda} - (n-1) S_{xx} \right)$$

Replace M by σ_{xx} & σ_{xy} by reusing the power-law



- No longer need to measure M independently \rightarrow removing uncertainty in M measurements
- All transport coefficients are from exactly the same region (Hall bar cross)

S_{yx} Sign Change



Open circuit voltage (J_x and $J_y=0$)

Nernst current:

$$J_y = \alpha_{yx} (-\nabla T)_x$$

Nernst current Hall current

$$S_{yx} = \left(\frac{\sigma_{xy}}{\sigma_{xx}} \right)' = \frac{1}{\sigma_{xx}} (\alpha_{yx} - \sigma_{yx} S_{xx})$$

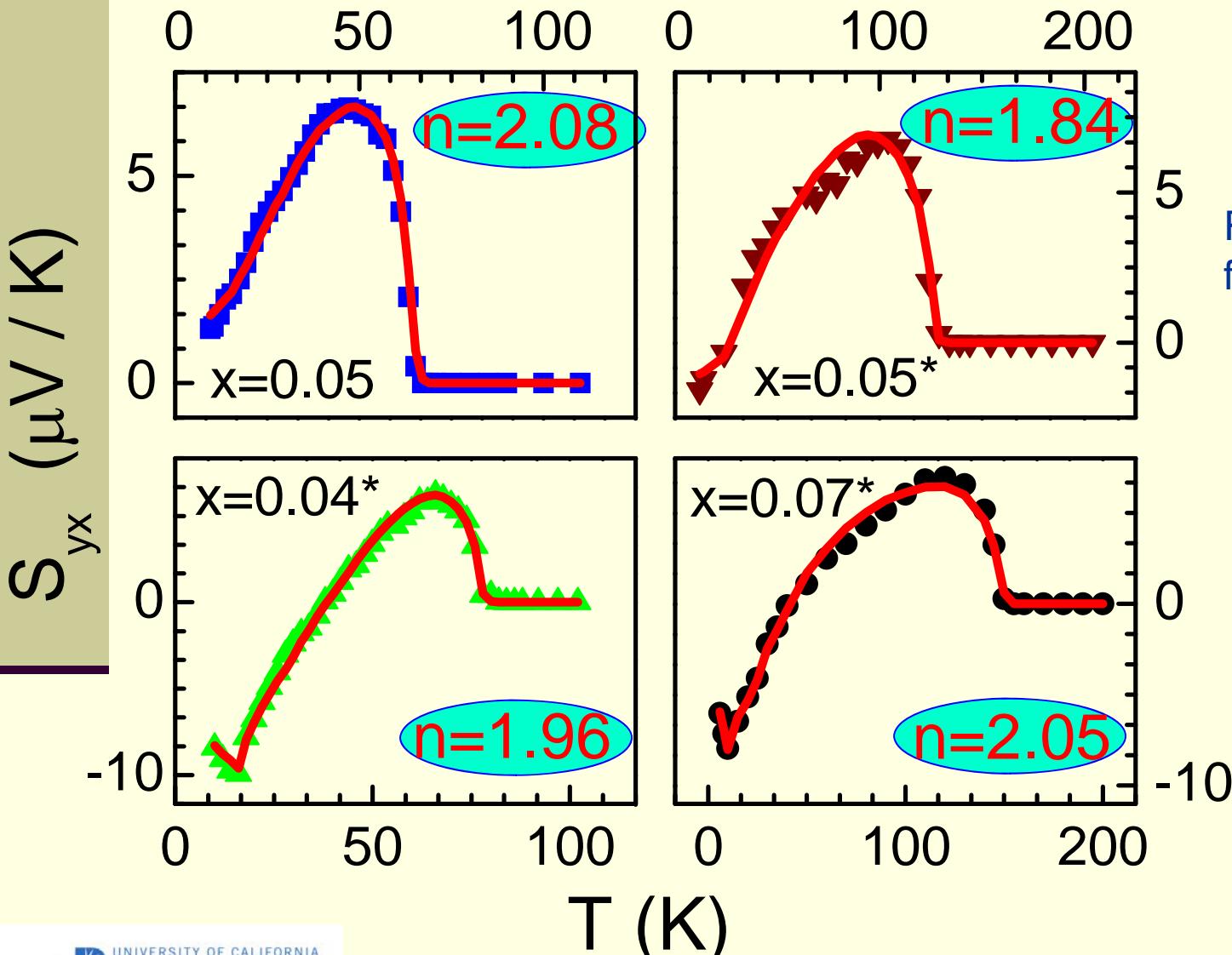
$$S_{yx} = \frac{\rho_{xy}}{\rho_{xx}} \left(T \frac{\pi^2 k_B^2}{3e} \frac{\lambda'}{\lambda} - (n-1) S_{xx} \right)$$

Nernst current Hall current

Sign change is only possible if $n > 1$, \rightarrow NOT skew scattering!

At low-T where S_{xx} is large, Hall current exceeds Nernst current \rightarrow Sign Change!

Exponent “n”



Red solid lines are fits to Mott relation

$$S_{yx} = \frac{\rho_{xy}}{\rho_{xx}} \left(T \frac{\pi^2 k_B^2 \lambda'}{3e \lambda} - (n-1) S_{xx} \right)$$

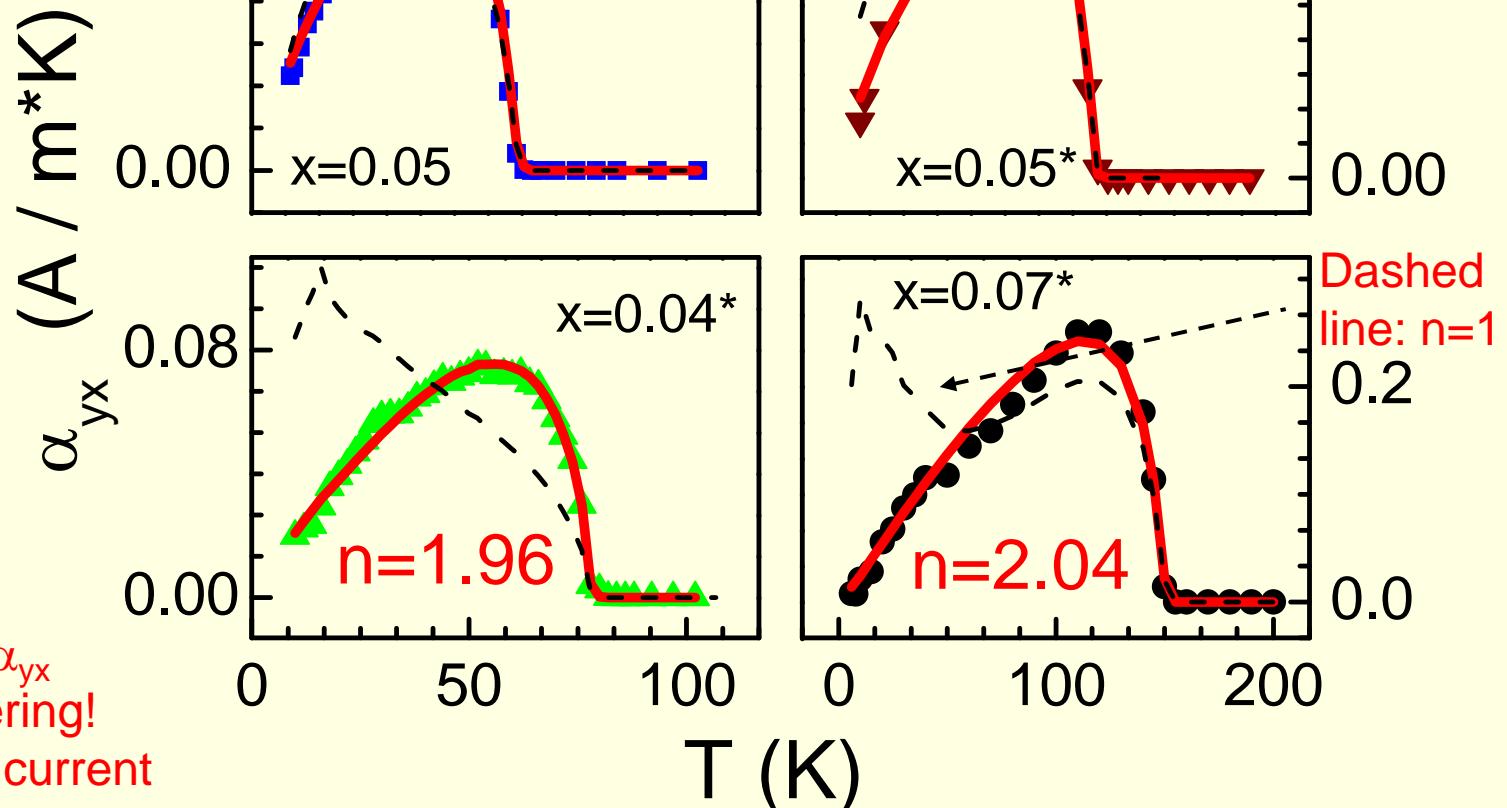
Mott relation works well for $n=2$!

Nernst Current: J_N

$$\alpha_{yx} = \frac{\rho_{xy}}{\rho_{xx}^2} \left(T \frac{\pi^2 k_B}{3e} \frac{\lambda}{\lambda} - (n-2) S_{xx} \right)$$

Nernst current:

$$J_N = -\alpha_{yx} (\nabla T)_x$$



$n=2 \rightarrow$ nothing in α_{yx}
depends on scattering!
→ intrinsic Nernst current

Intrinsic vs. Side Jump

- $n=2 \rightarrow$ intrinsic or side jump (SJ)

In GaMnAs,

- ❖ Intrinsic mechanism can account for most of AHE magnitude
- ❖ SJ displacement $\Delta y < 0.1$ nm, and MFP is about ~ 20 nm.
Hall angle $\Theta_H < 0.05$ for SJ; observed Hall angle $\Theta_H \sim 0.1$
- ❖ AHE/ANE is likely dominated by intrinsic mechanism

Mott relation holds for AHE/ANE with intrinsic (Berry's phase) mechanism.

Meaning of λ/λ'

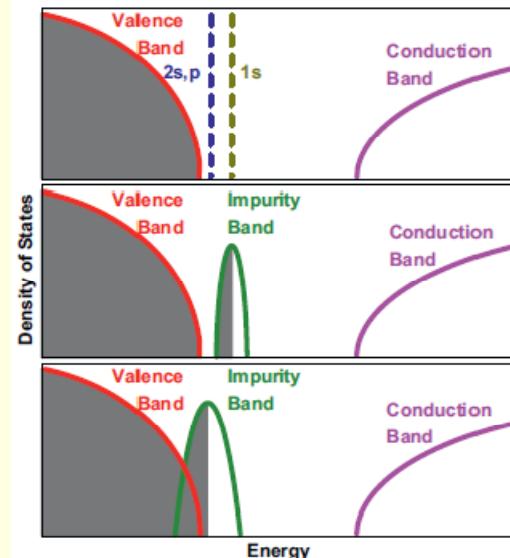
- λ depends on Fermi energy
- We obtain the value of λ/λ'

	#4*	#5	#5*	#7*
n	1.96	1.95	1.8	2.04
$\lambda/\lambda'(\text{eV})$	0.046	0.049	0.040	0.096

For small ε_F , we assume $\lambda \propto \varepsilon_F^\alpha$

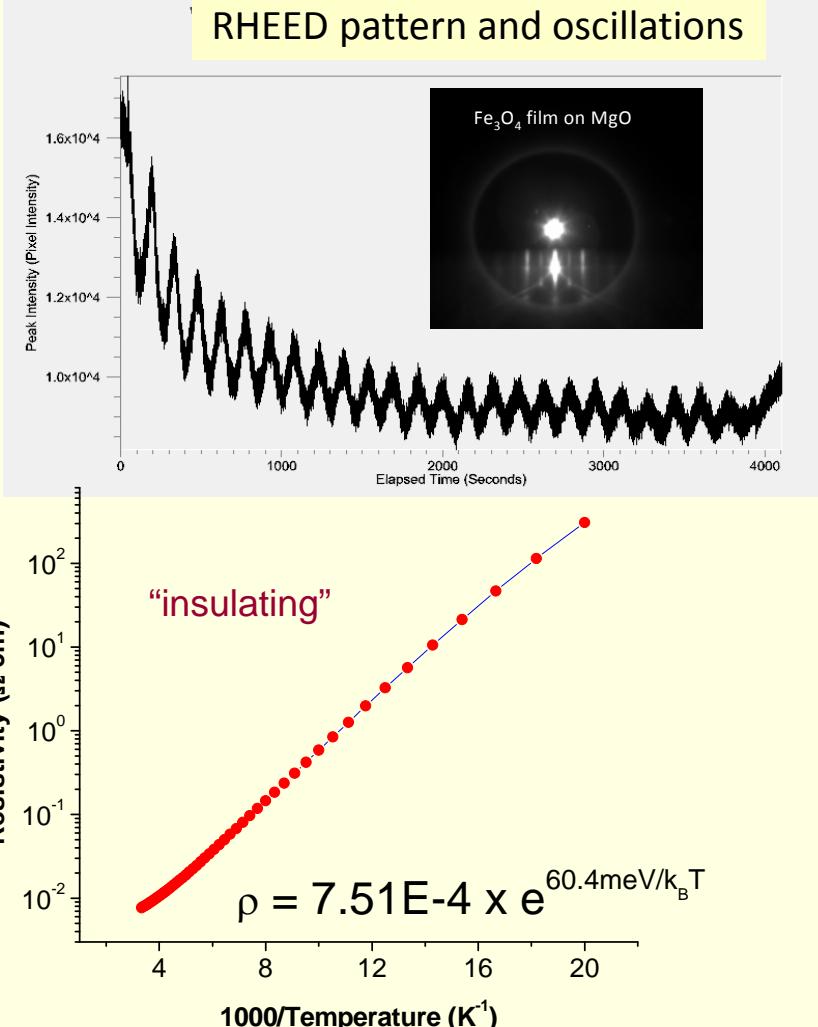
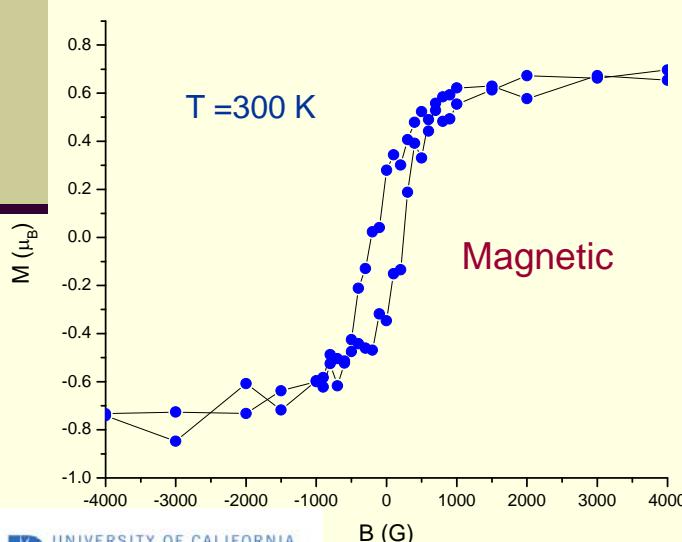
$$\frac{\lambda}{\lambda'} = \frac{\varepsilon_F}{\alpha} \approx 0.05 - 0.1 \text{ eV}$$

For $p \sim 10^{19} \text{ cm}^{-3}$, obtained ε_F is an order smaller than expected from valence band, suggesting the impurity band picture.

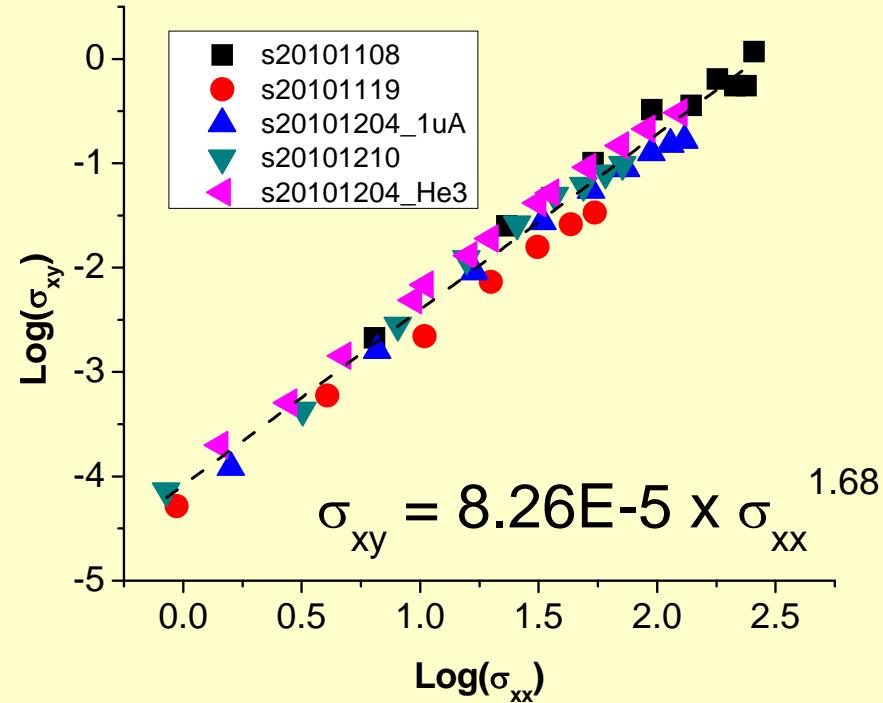
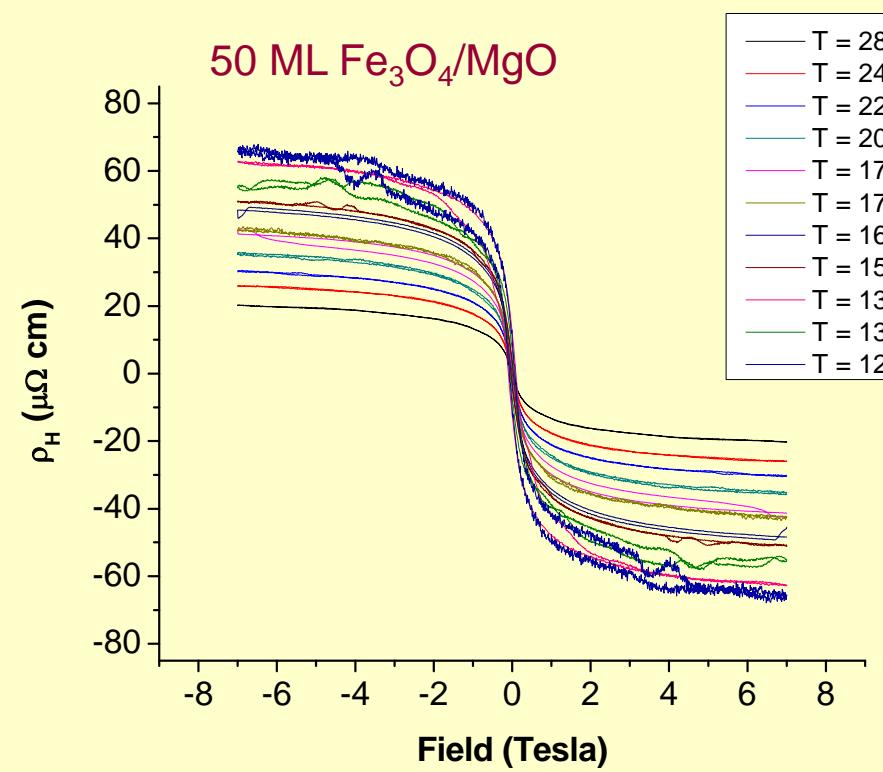


Magnetic Insulator Films: $\text{Fe}_3\text{O}_4/\text{MgO}$ (001)

- Fe_3O_4 films are epitaxially grown with laser MBE



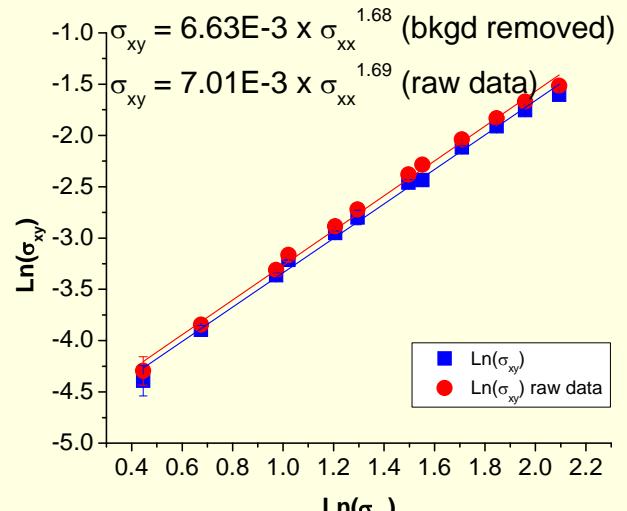
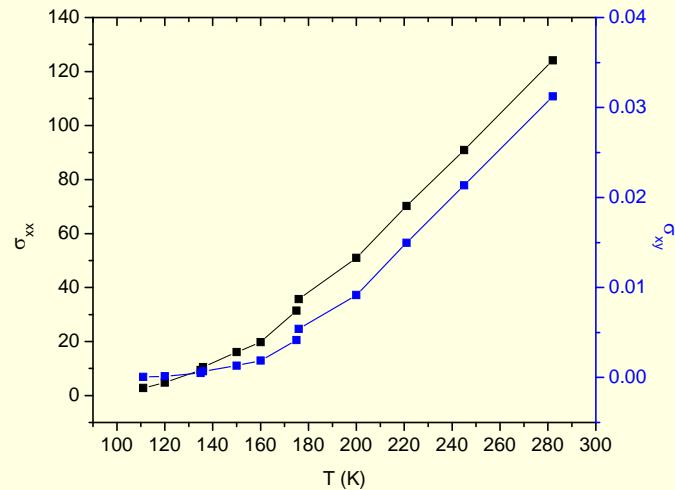
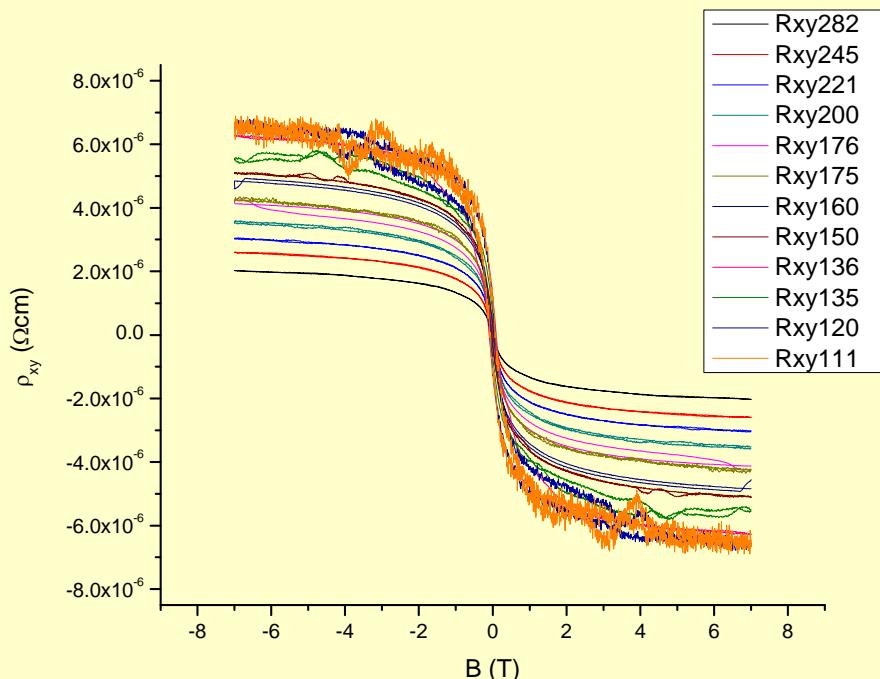
Power-Law Scaling



AHE signal above Verwey temperature

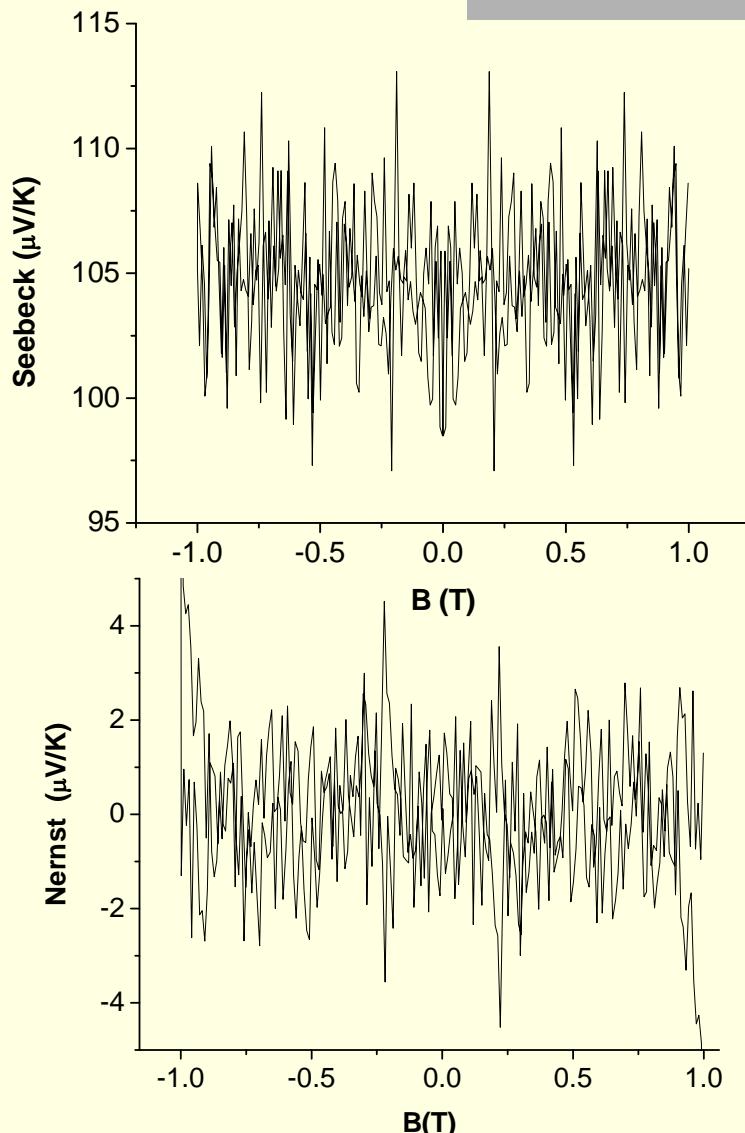
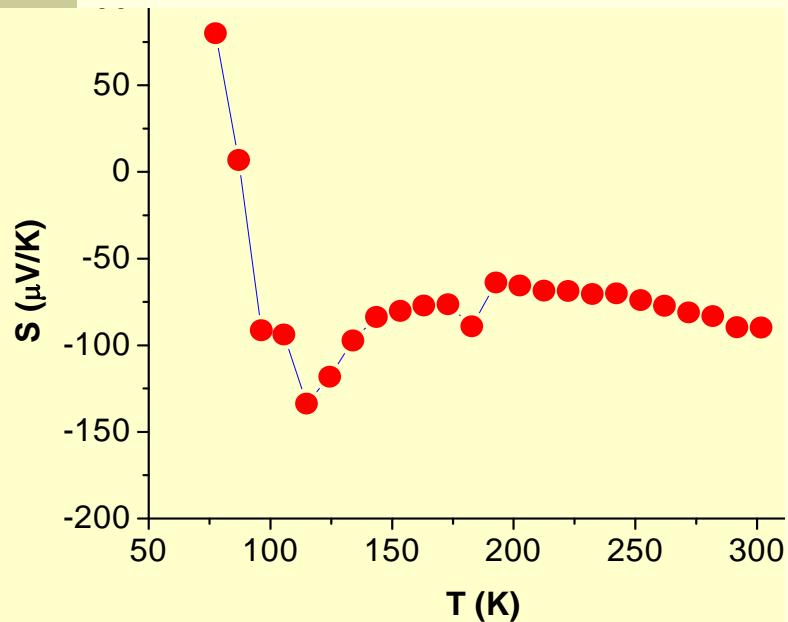
A much smaller $n=0.32$!
Consistent with other reported value.

AHE and Power-Law



Thermoelectric Measurements

Zero-field Seebeck

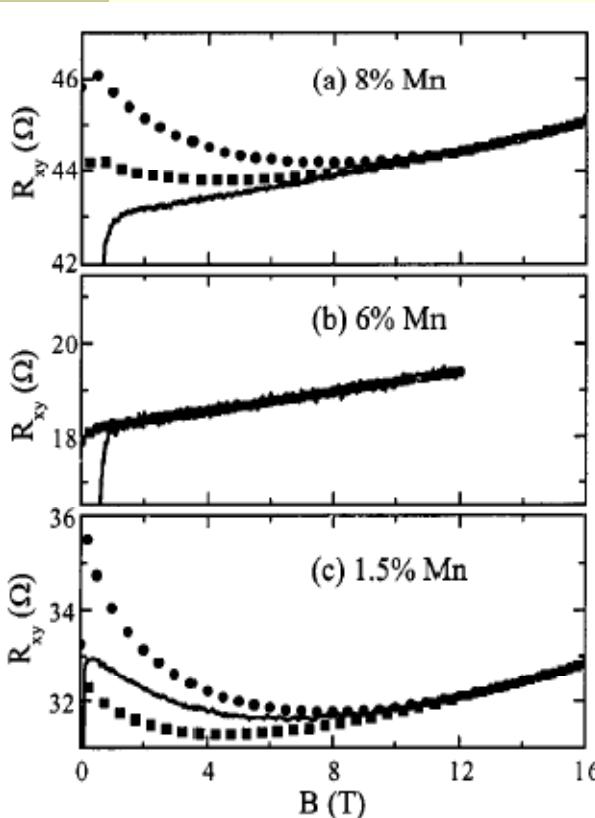


Summary

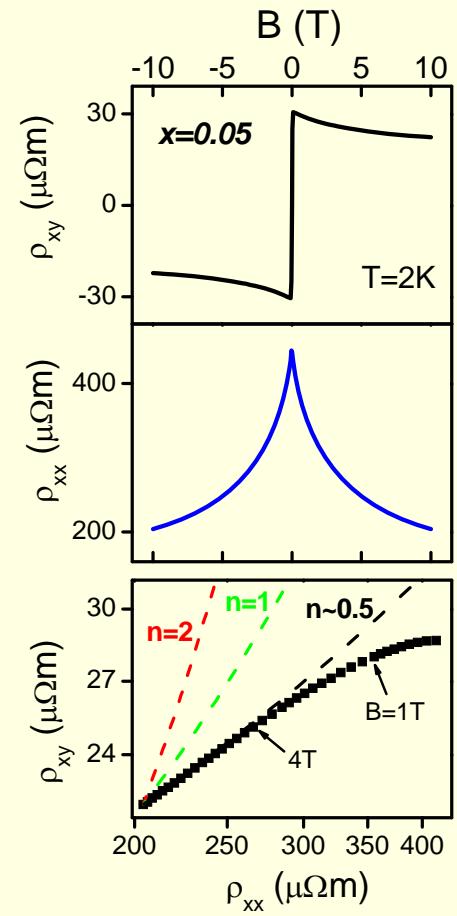
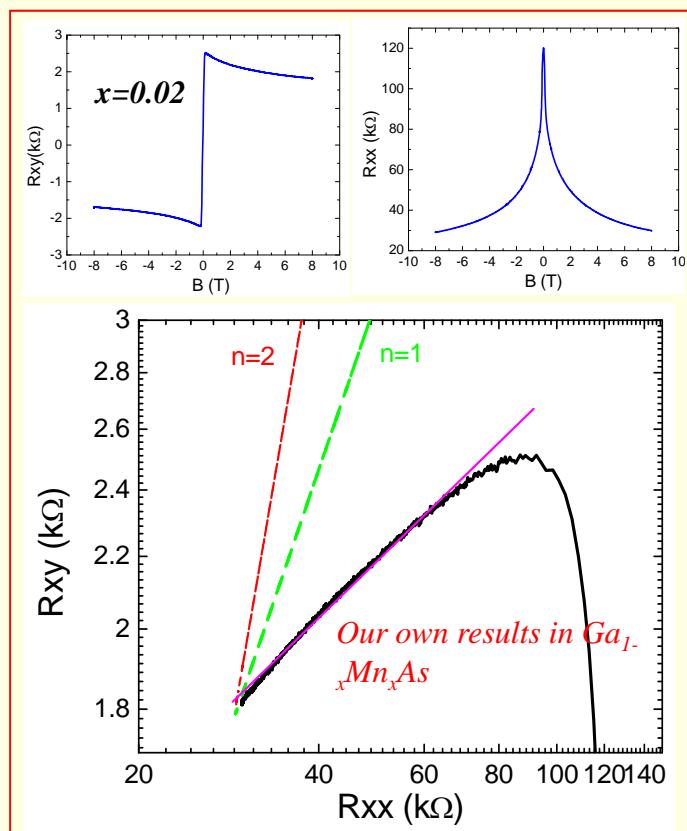
- Large ANE is observed in GaMnAs [$\lambda=\lambda(\varepsilon_F)$].
- AHE and ANE share the same physical origin.
- “n=2” is obtained from zero-field AHE and ANE without uncertainty in magnetization measurements.
- Our results suggest intrinsic Nernst current J_N .
- Mott relation is experimentally validated for scattering rate-independent anomalous transport.
- Small magnitude of λ/λ' suggests impurity band picture.
- “n=0.3” is found in epitaxial Fe_3O_4 films, but ANE has not been observed.

Discussion: 1. Effect of High B-Field

In most experiments with in-plane anisotropy samples, high magnetic fields are used, but the effect of B-field on the power-law scaling is unclear. Which field is the proper one?

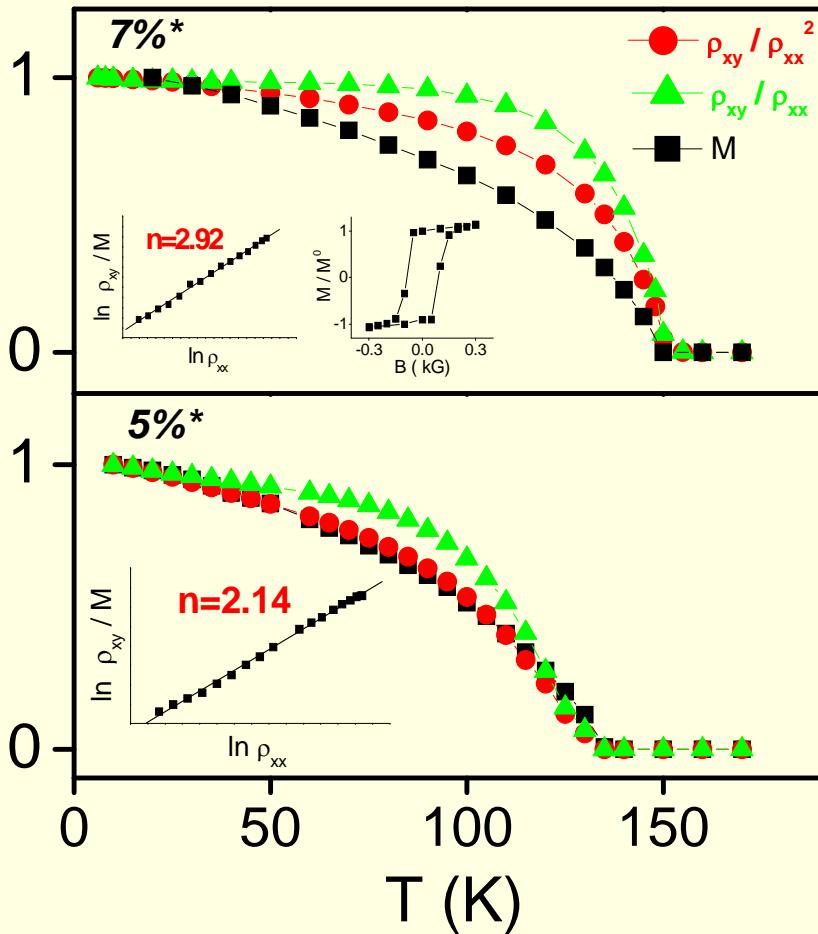


Edmonds et al. JAP (03)



Find Exponent n

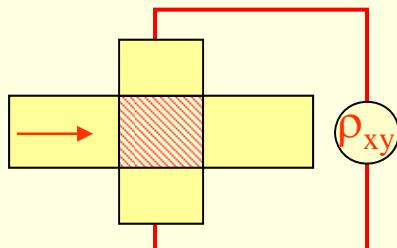
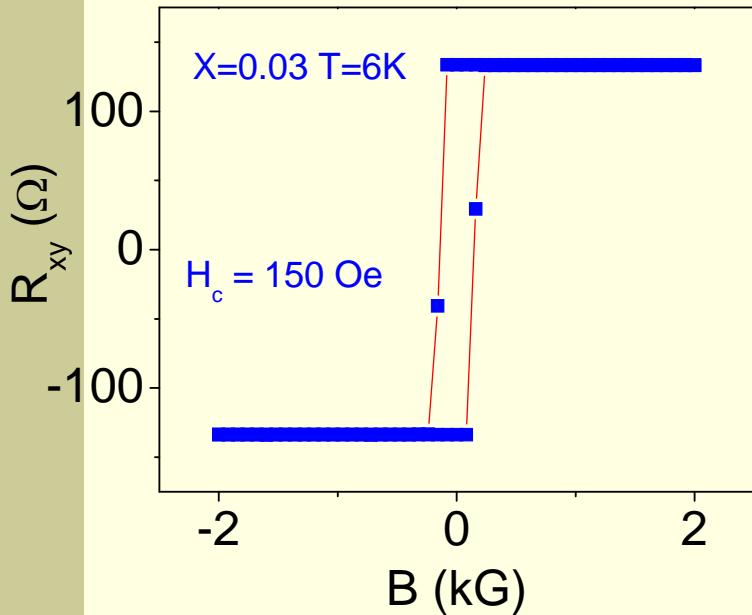
A. U.



- Exponent is not always one or two
- Separate magnetization measurement is needed
- Hall loops are different from MH loops

R_{xy} and M Measurements

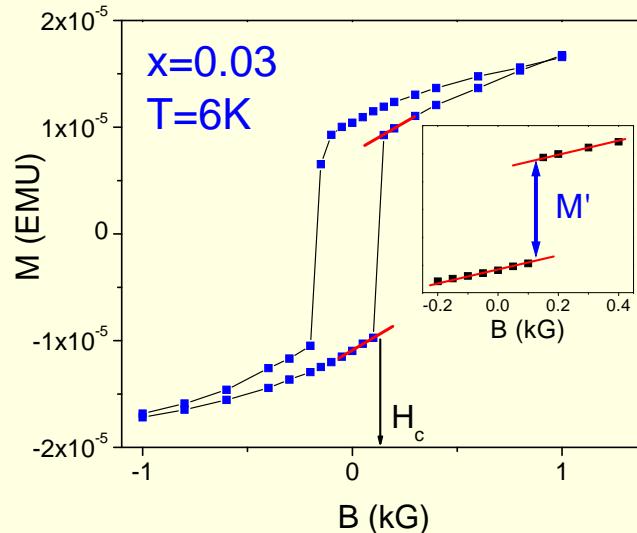
Transport



Probing area $\sim (100 \mu\text{m})^2$

Good

Magnetization



Bad

