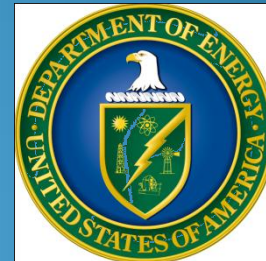


# Hofstadter's Fractal Energy Spectrum in Twisted Bilayer Graphene

Feng Liu<sup>1</sup>, Zhengfei Wang<sup>1,2</sup>, Mei-Yin Chou<sup>2</sup>

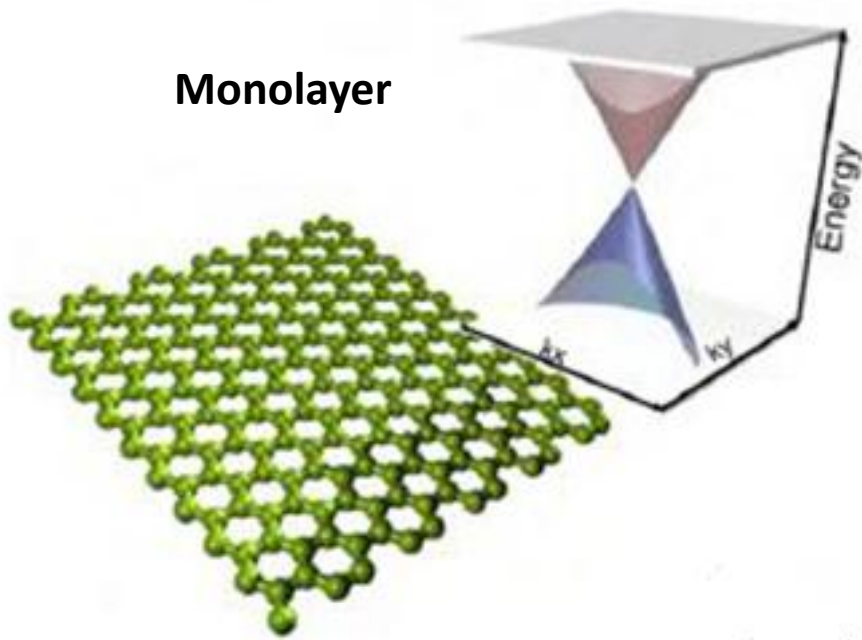
<sup>1</sup>*Department of Materials Science and Engineering, University of Utah*

<sup>2</sup>*School of Physics, Georgia Institute of Technology*

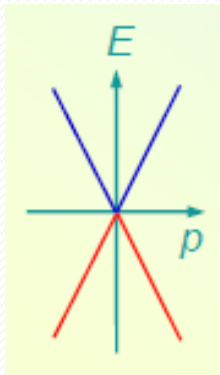
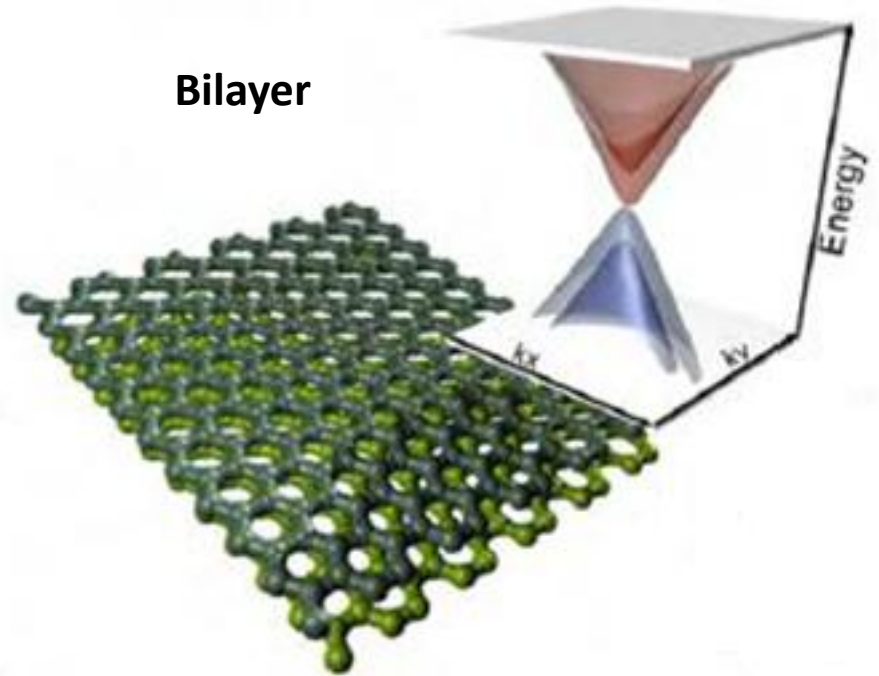


# Band Structure of Monolayer and Bilayer Graphene

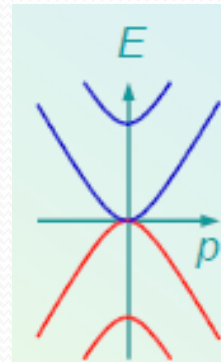
Monolayer



Bilayer

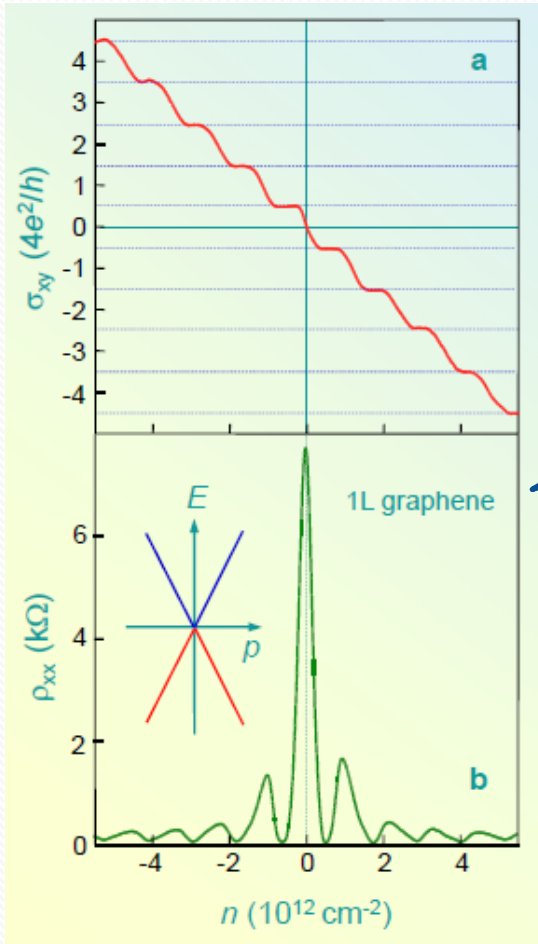


$$E = \pm \hbar v_f k$$



$$E = \pm \frac{\gamma_1}{2} \pm \sqrt{(\hbar v_f k)^2 + \left(\frac{\gamma_1}{2}\right)^2}$$

# Landau Levels of Monolayer and Bilayer Graphene



$$\varepsilon_N = \pm v_F \sqrt{2e\hbar BN}$$

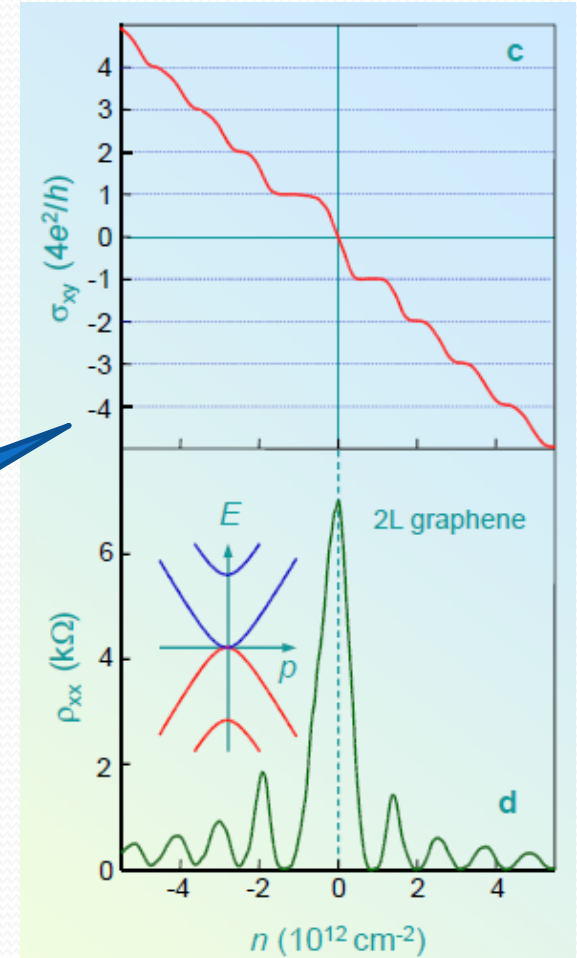
$$v_F = (\sqrt{3}/2)\gamma_0 a/\hbar$$

$$\sigma_{xy} = \pm 4\left(N + \frac{1}{2}\right) \frac{e^2}{h}$$

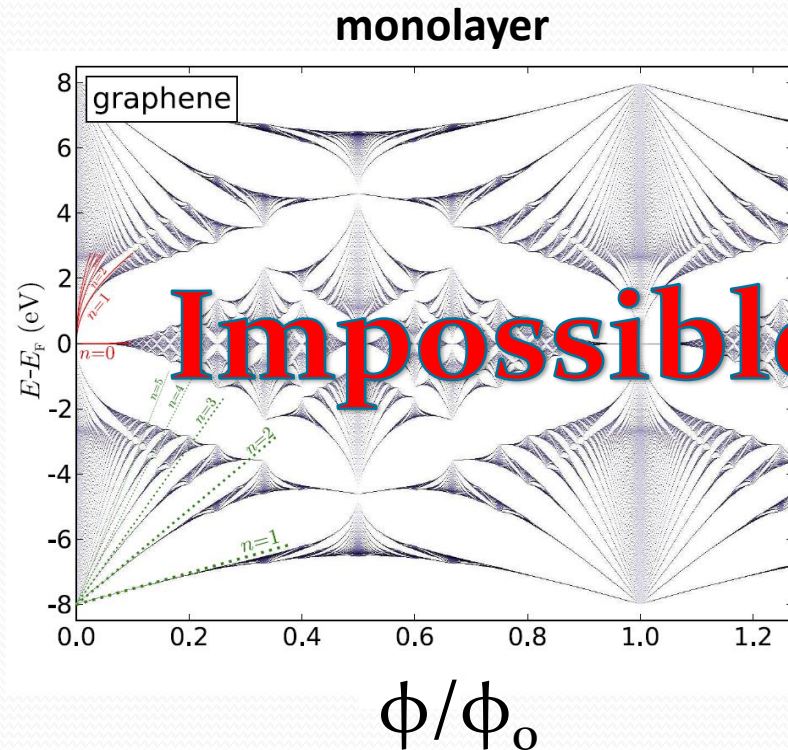
$$\varepsilon_N = \pm \hbar \omega_c \sqrt{N(N-1)}$$

$$\omega_c = eB/m$$

$$\sigma_{xy} = \pm 4N \frac{e^2}{h}$$

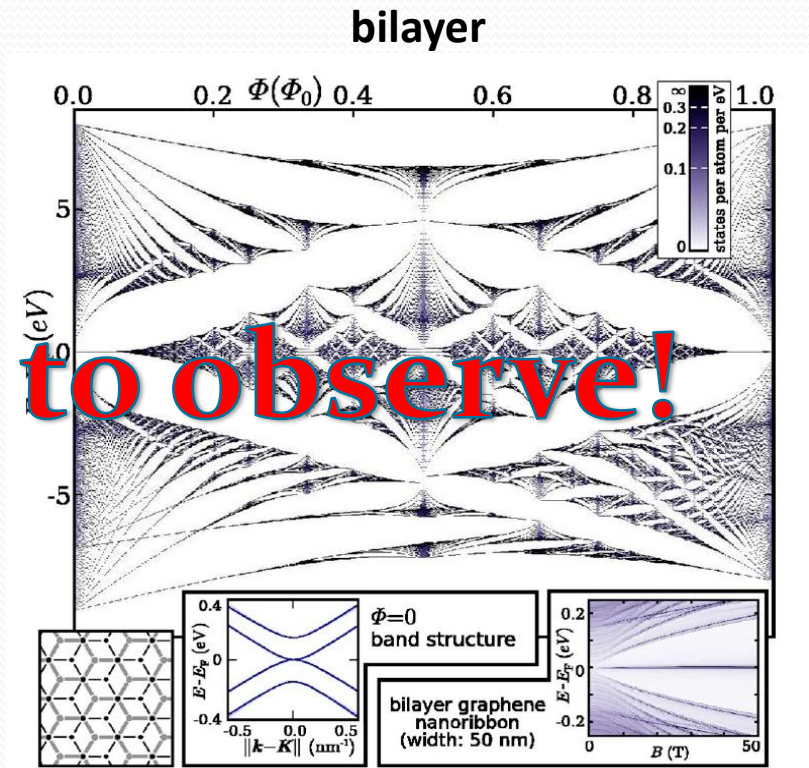


# Hofstadter Butterflies of Monolayer and Bilayer Graphene



N. Nemec *et al.*

Phys. Rev. B 74, 165411 (2006)



N. Nemec *et al.*

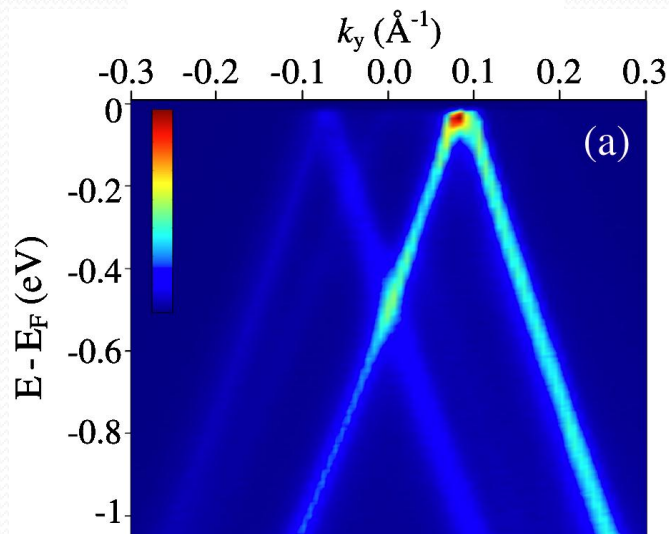
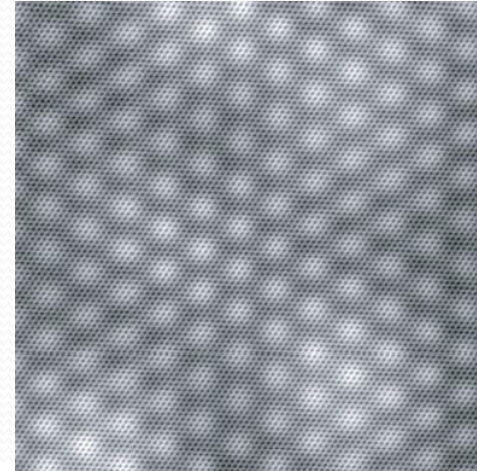
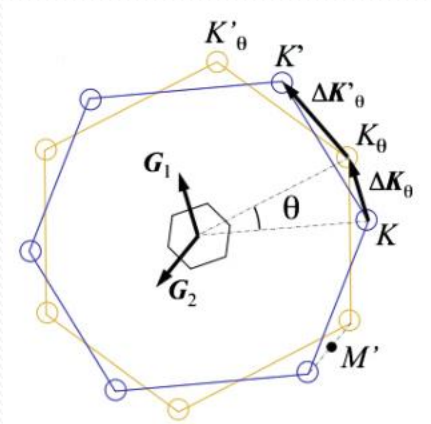
Phys. Rev. B 75, 201404 (R) (2007)

$$S = \frac{3\sqrt{3}}{2} a^2$$

$$\phi/\phi_0 = 1, \quad B = 79098 \text{ T}$$

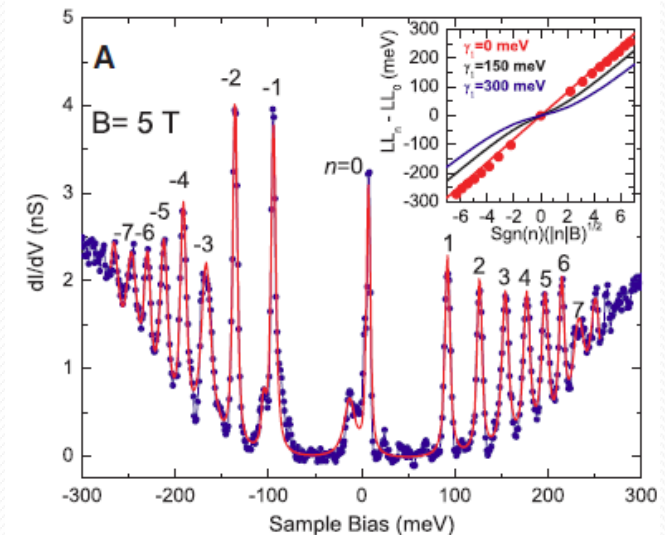


# Decoupling Behavior of Multilayer Epitaxial Graphene



M. Sprinkle *et al.*

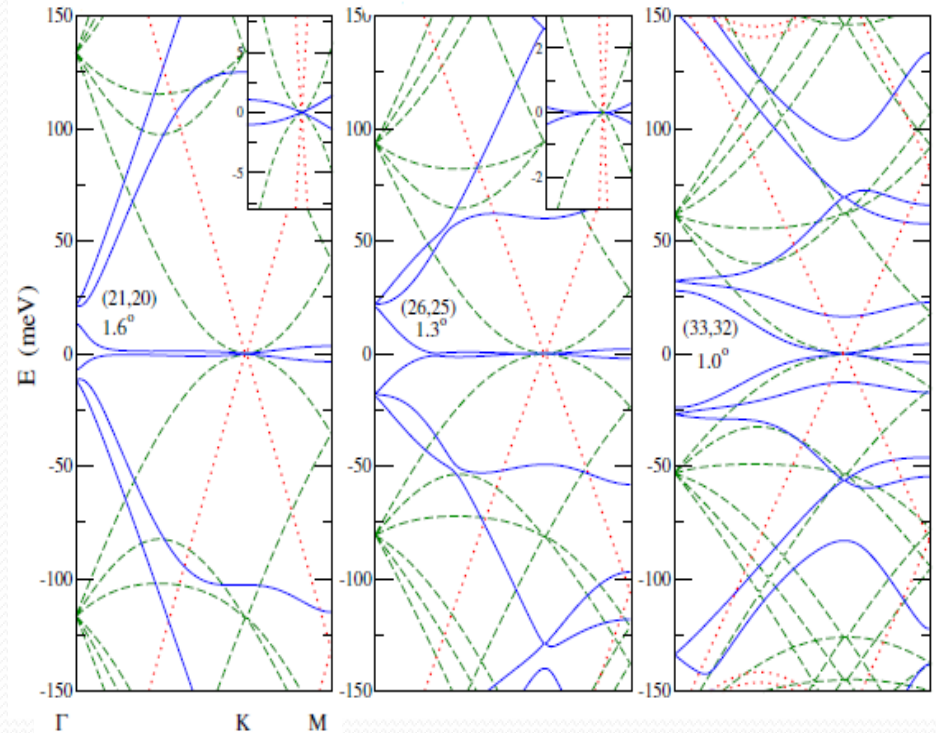
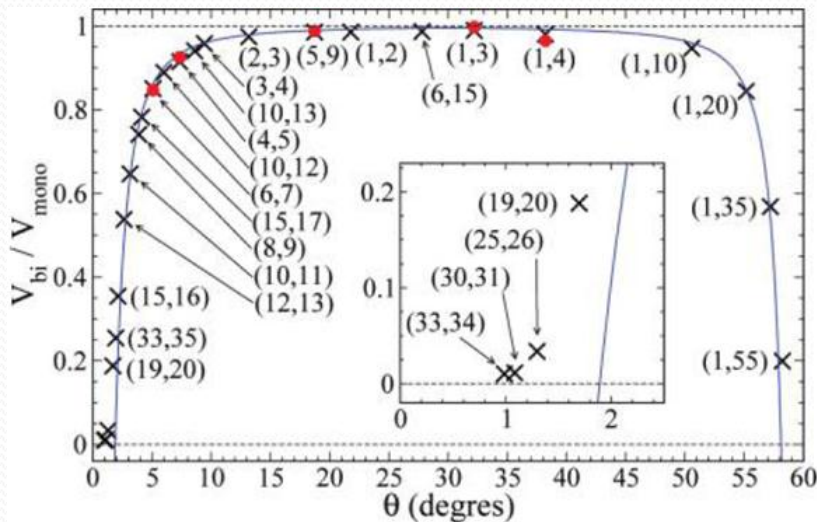
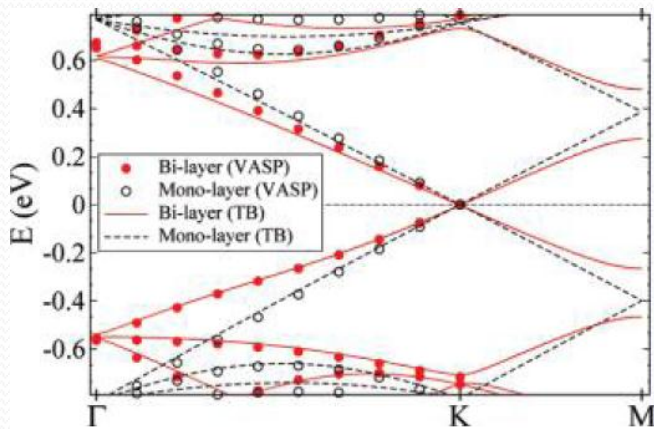
Phys. Rev. Lett. 103, 226803 (2009)



D. L. Miller *et al.*

Science 324, 924 (2009)

# Band Structure of Twisted Bilayer Graphene



E. Suárez Morell *et al.*

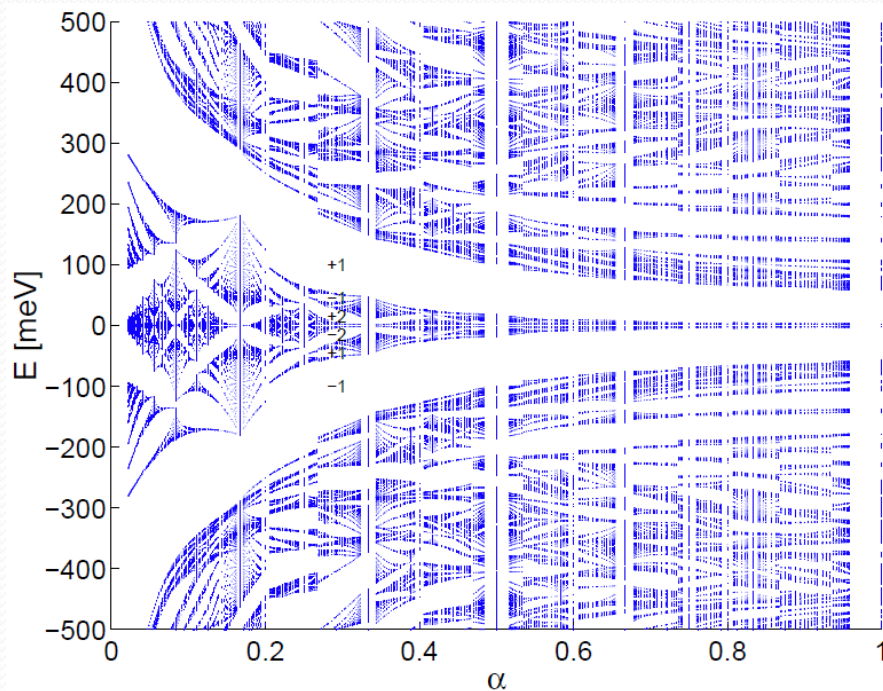
Phys. Rev. B 82, 121407(R) (2010)

G. T. de Laissardiere *et al.*

Nano Lett. 10, 804 (2010)

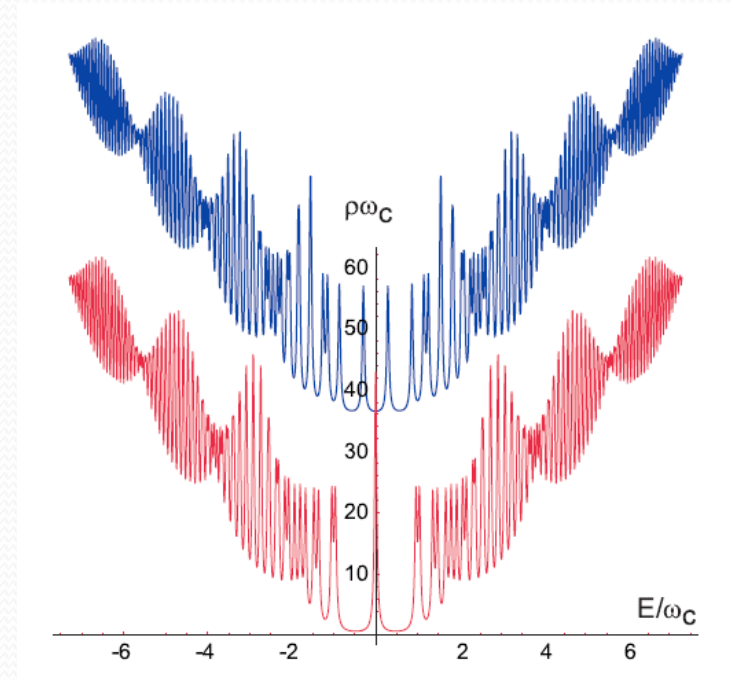
# Recent Results of Twisted Bilayer Graphene in “Low” Magnetic Field

## Morie Butterfly



R. Bistritzer and A.H. MacDonald  
arXiv:1101.2606v1

## Dirac Comb



M. Kindermann and E.J. Mele  
arXiv:1106.0204v1

# Motivation

- How do the electronic properties of twisted bilayer graphene change with the twist angle?
- How do the LLs evolve as the twist angle changes?
- Will there be a difference between twisted bilayer graphene with commensurate and incommensurate angles?



# Lanczos Recursive Method

Real space Hamiltonian (Hermitian matrix)

$$H = \sum_{i,j} t_{ij} \exp\left(\frac{ie}{\hbar} \int_i^j \vec{A} \cdot d\vec{l}\right) a_i^+ a_j \quad \vec{A} = (0, Bx)$$

Construct a new orthogonal basis

$$\{|\Phi_0\rangle, |\Phi_1\rangle, |\Phi_2\rangle, \dots, |\Phi_N\rangle\} \quad \begin{array}{l} |\Phi_0\rangle \text{ is the initial state localized at one atom site} \\ N \text{ is the total number of atoms} \end{array}$$

$$|\Phi_0\rangle$$

$$a_0 = \langle \Phi_0 | H | \Phi_0 \rangle$$

$$|\tilde{\Phi}_1\rangle = b_1 |\Phi_1\rangle = H |\Phi_0\rangle - a_0 |\Phi_0\rangle$$

$$a_1 = \langle \Phi_1 | H | \Phi_1 \rangle \quad b_1 = \sqrt{\langle \tilde{\Phi}_1 | \tilde{\Phi}_1 \rangle}$$

$$|\tilde{\Phi}_2\rangle = b_2 |\Phi_2\rangle = H |\Phi_1\rangle - a_1 |\Phi_1\rangle - b_1 |\Phi_0\rangle$$

$$a_1 = \langle \Phi_2 | H | \Phi_2 \rangle \quad b_2 = \sqrt{\langle \tilde{\Phi}_1 | \tilde{\Phi}_2 \rangle}$$

.....

## Recursive relation

$$|\tilde{\Phi}_{N+1}\rangle = b_{N+1} |\Phi_{N+1}\rangle = H |\Phi_N\rangle - a_N |\Phi_N\rangle - b_N |\Phi_{N-1}\rangle$$

$$a_N = \langle \Phi_N | H | \Phi_N \rangle \quad b_N = \sqrt{\langle \tilde{\Phi}_N | \tilde{\Phi}_N \rangle}$$

Hamiltonian in the new basis

$$H_{ij} = \langle \Phi_i | H | \Phi_j \rangle \quad H = \begin{pmatrix} a_0 & b_1 & 0 & \cdots \\ b_1 & a_1 & b_2 & \cdots \\ 0 & b_2 & a_2 & \cdots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$

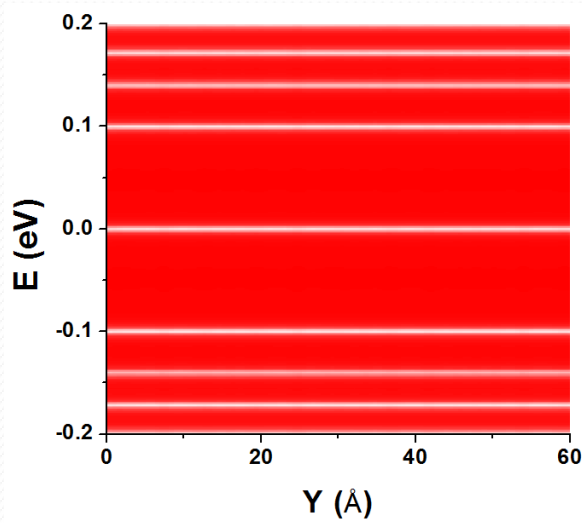
Real space Green's function at the initial state (continued fraction expansion)

$$\langle \Phi_0 | G^r(E + i\eta) | \Phi_0 \rangle = \frac{1}{E + i\eta - a_0 - \frac{b_1^2}{E + i\eta - a_1 - \frac{b_2^2}{E + i\eta - a_2 - \cdots}}}$$

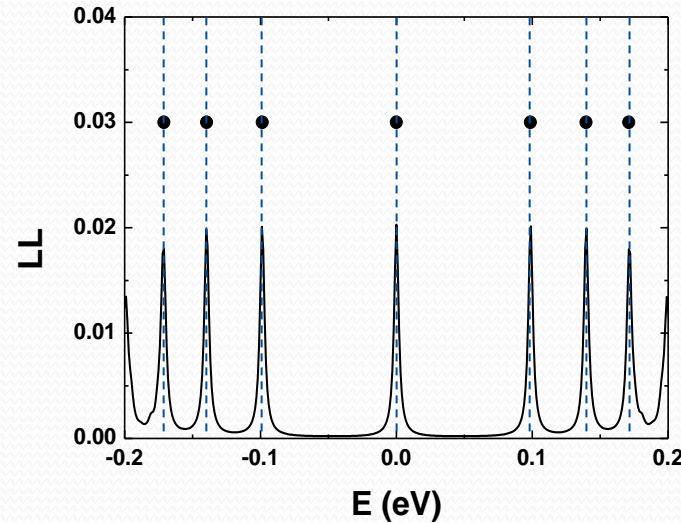
$$LDOS = -\frac{1}{\pi} \text{Im} G^r(E + i\eta) \quad \cdots$$

**140nm×140nm, over 1.5million atoms!**

# Landau Levels of Monolayer and Bilayer Graphene

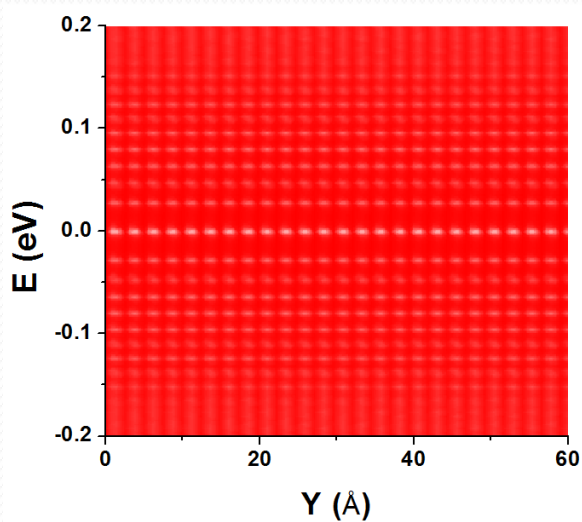


Monolayer graphene

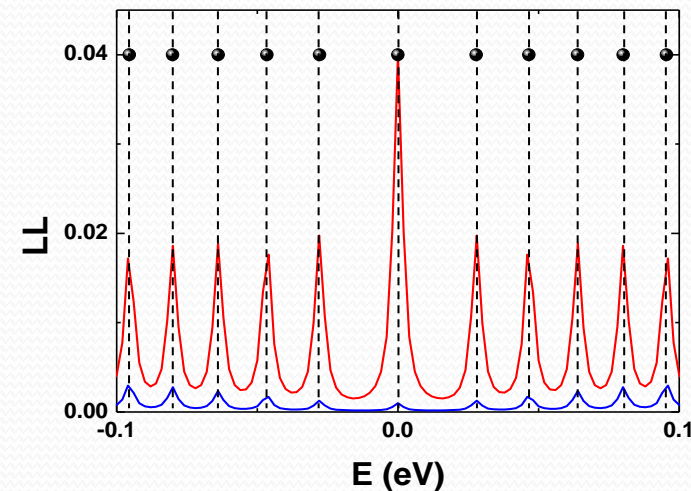


$$E_n = \text{sgn}(n)\Delta\sqrt{|n|}$$

$$\Delta = \sqrt{2eBv_F^2\hbar}$$

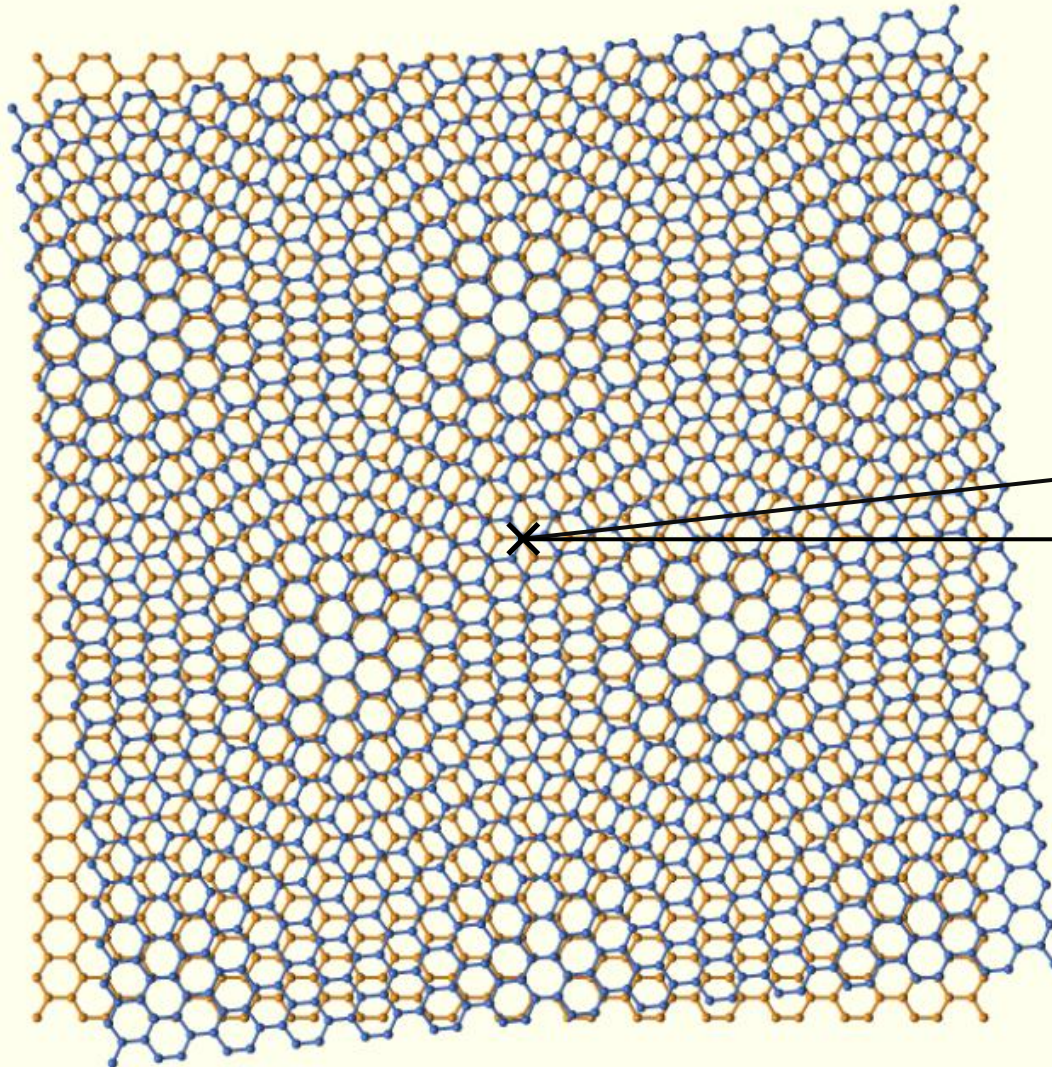


AB bilayer graphene



$$E_n = \frac{\text{sgn}(n)}{\sqrt{2}} \left[ (2|n|+1)\Delta^2 + \gamma_1^2 - \sqrt{\gamma_1^4 + 2(2|n|+1)\Delta^2\gamma_1^2 + \Delta^4} \right]^{1/2}$$

# Twisted Bilayer Graphene



Starting from AB stacking bilayer graphene, bottom layer is fixed and top layer is twisted.

Twist center is at A atom in the top layer with a neighboring B atom in the bottom layer.

$\theta$  Commensurate angle

$$\theta = \cos^{-1} \left( \frac{3q^2 - p^2}{3q^2 + p^2} \right)$$

**S. Shallcross *et al.***  
**Phys. Rev. B 81, 1 (2010)**

Other  $\theta$ 's are  
incommensurate angles



# Tight-binding Parameterization

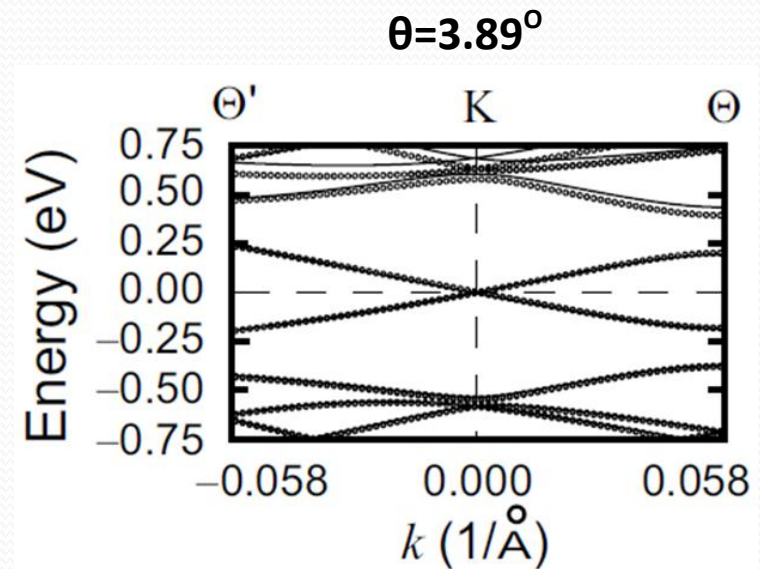
$$H = \sum_{\mu, \nu=1,2} \sum_{l,j} t_{\mu l, \nu j} \exp\left[\frac{ie}{\hbar} \int_{\vec{r}_{\mu l}}^{\vec{r}_{\nu j}} \vec{A} \cdot d\vec{l}\right] |\mu l\rangle \langle \nu j|$$

TB parameters are obtained by fitting the TB bands to reproduce first-principles band structures

$$t_{\mu l, \nu j} = \begin{cases} \gamma_1 \exp[\lambda_1 (1 - |\vec{r}_{\mu l} - \vec{r}_{\nu j}|/a)] & (\mu = \nu) \\ \gamma_2 \exp[\lambda_2 (1 - |\vec{r}_{\mu l} - \vec{r}_{\nu j}|/c)] & (\mu \neq \nu) \end{cases}$$

$a = 1.42 \text{ \AA}$	$\gamma_1 = -2.7 \text{ eV}$	$\lambda_1 = 3.15$
$c = 3.35 \text{ \AA}$	$\gamma_2 = 0.48 \text{ eV}$	$\lambda_2 = 7.42$

G. T. de Laissardiere *et al.*  
Nano Lett. 10, 804 (2010)



# Landau Levels as a Function of Position

Commensurate , B=10T

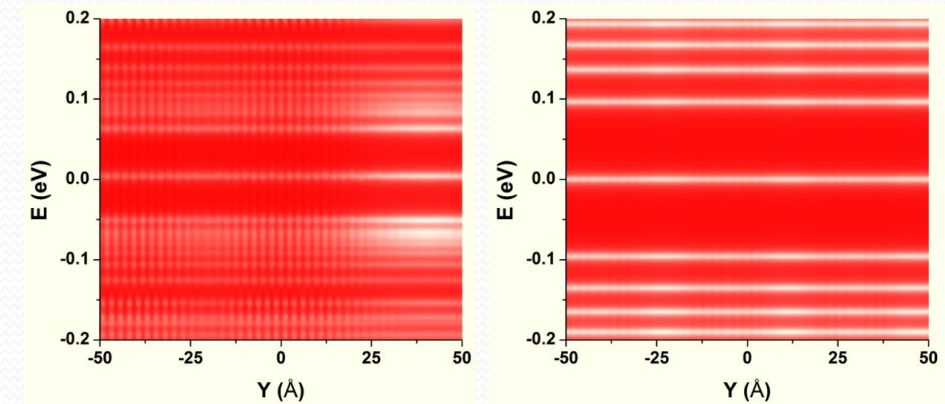
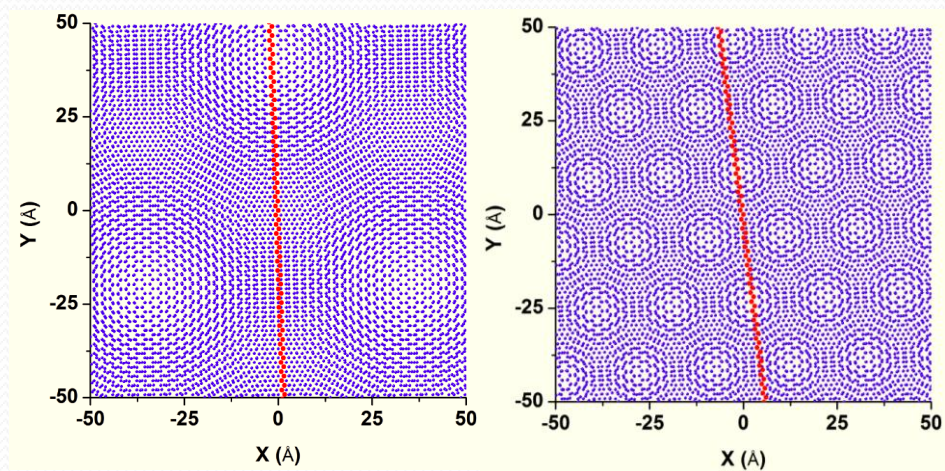
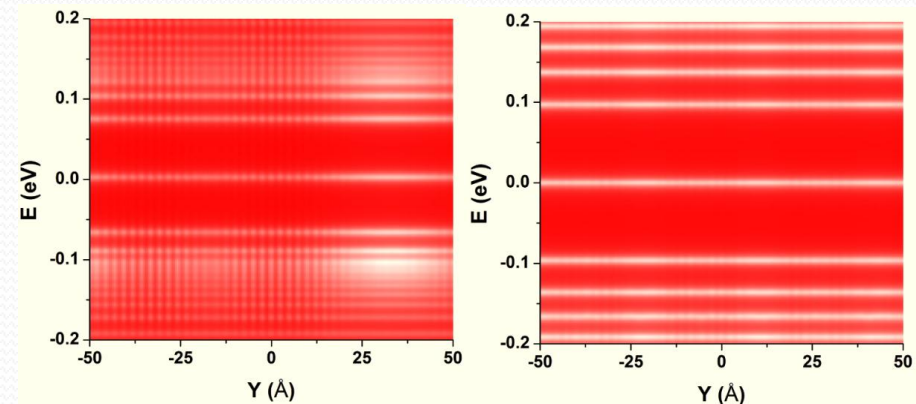
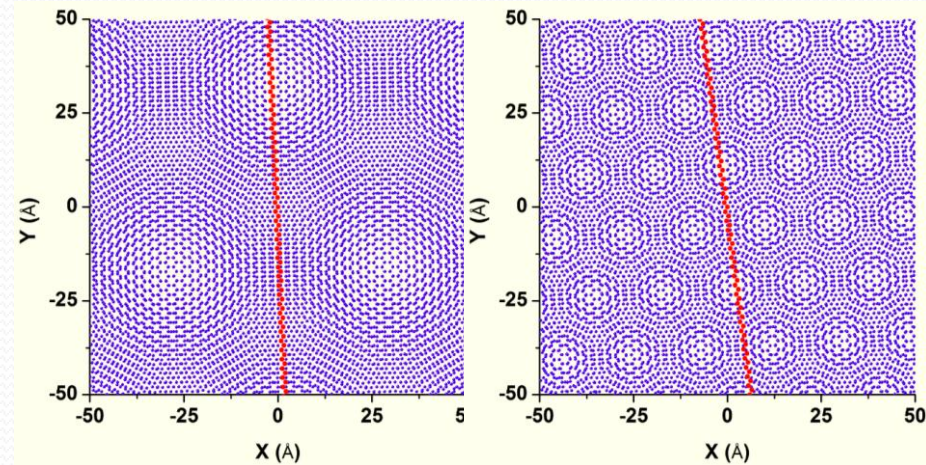
Incommensurate , B=10T

$\theta=2.4718^\circ$

$\theta=7.56507^\circ$

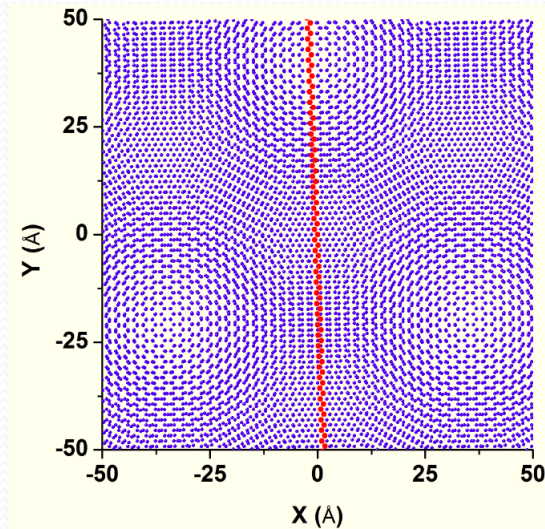
$\theta=2.0^\circ$

$\theta=7.0^\circ$

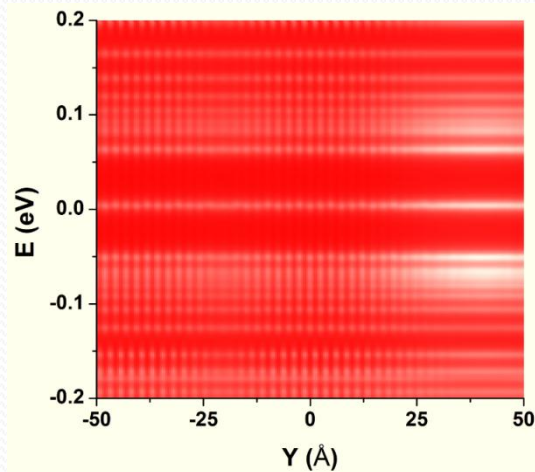


# Landau Levels as a Function of Position

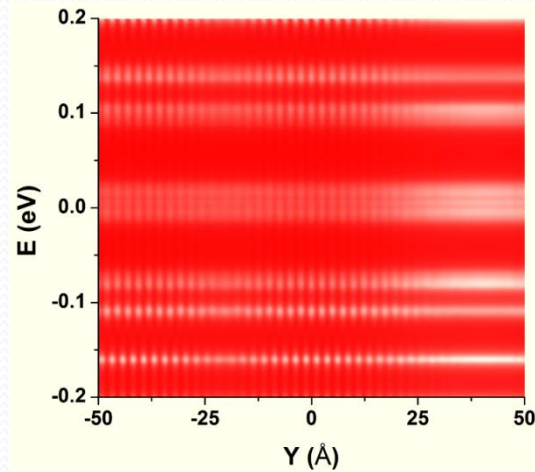
Incommensurate,  $\theta=2.0^\circ$



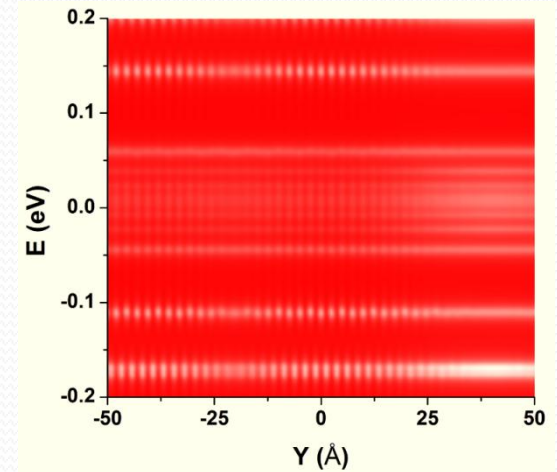
**B=10T**



**B=30T**



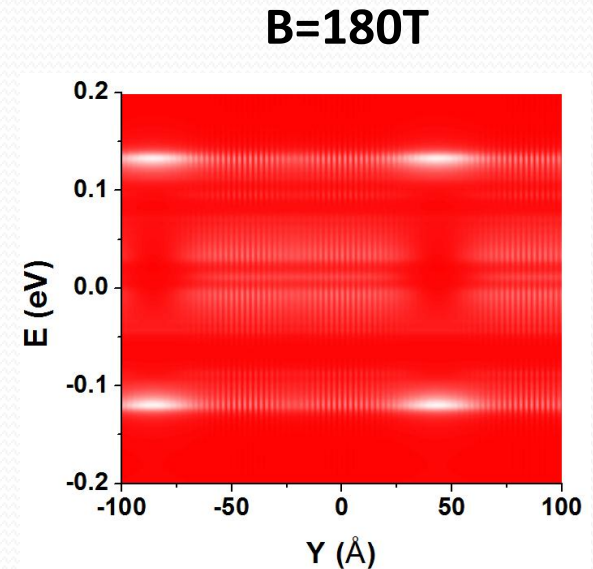
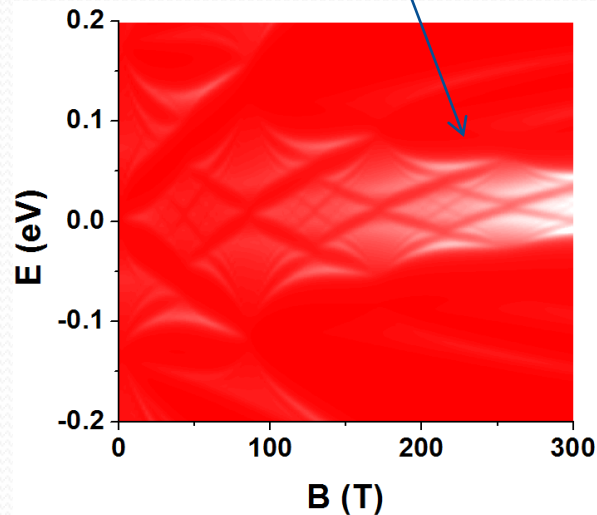
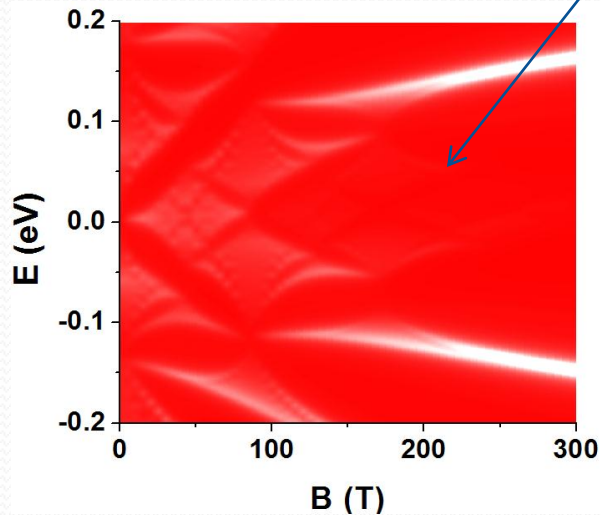
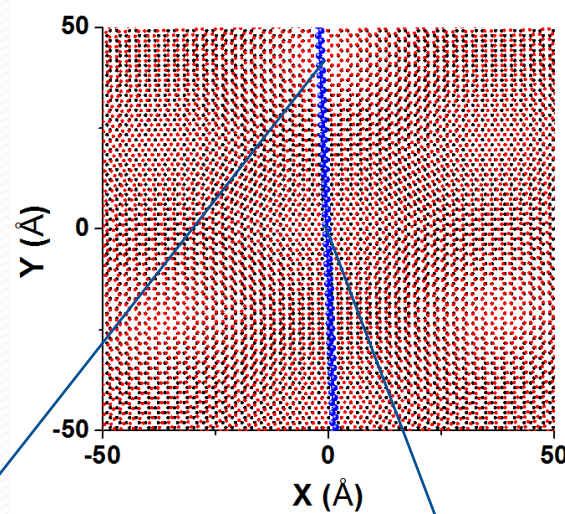
**B=60T**





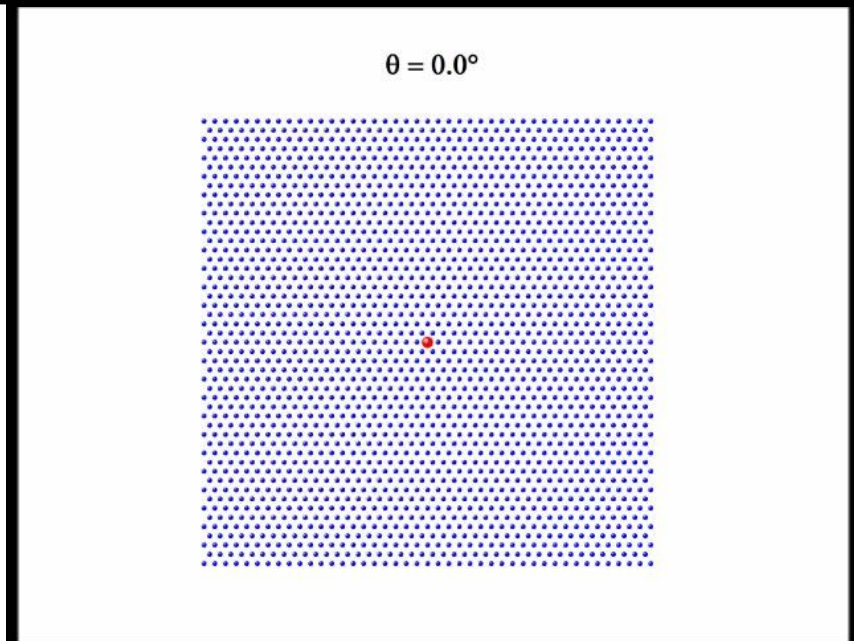
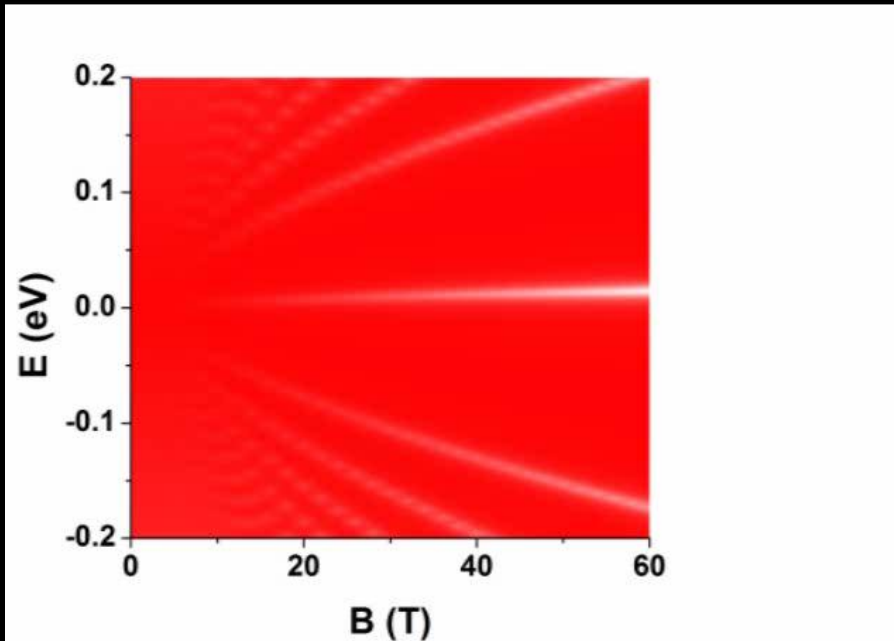
# Landau Levels as a Function of Position

commensurate,  $\theta=1.8901^\circ$

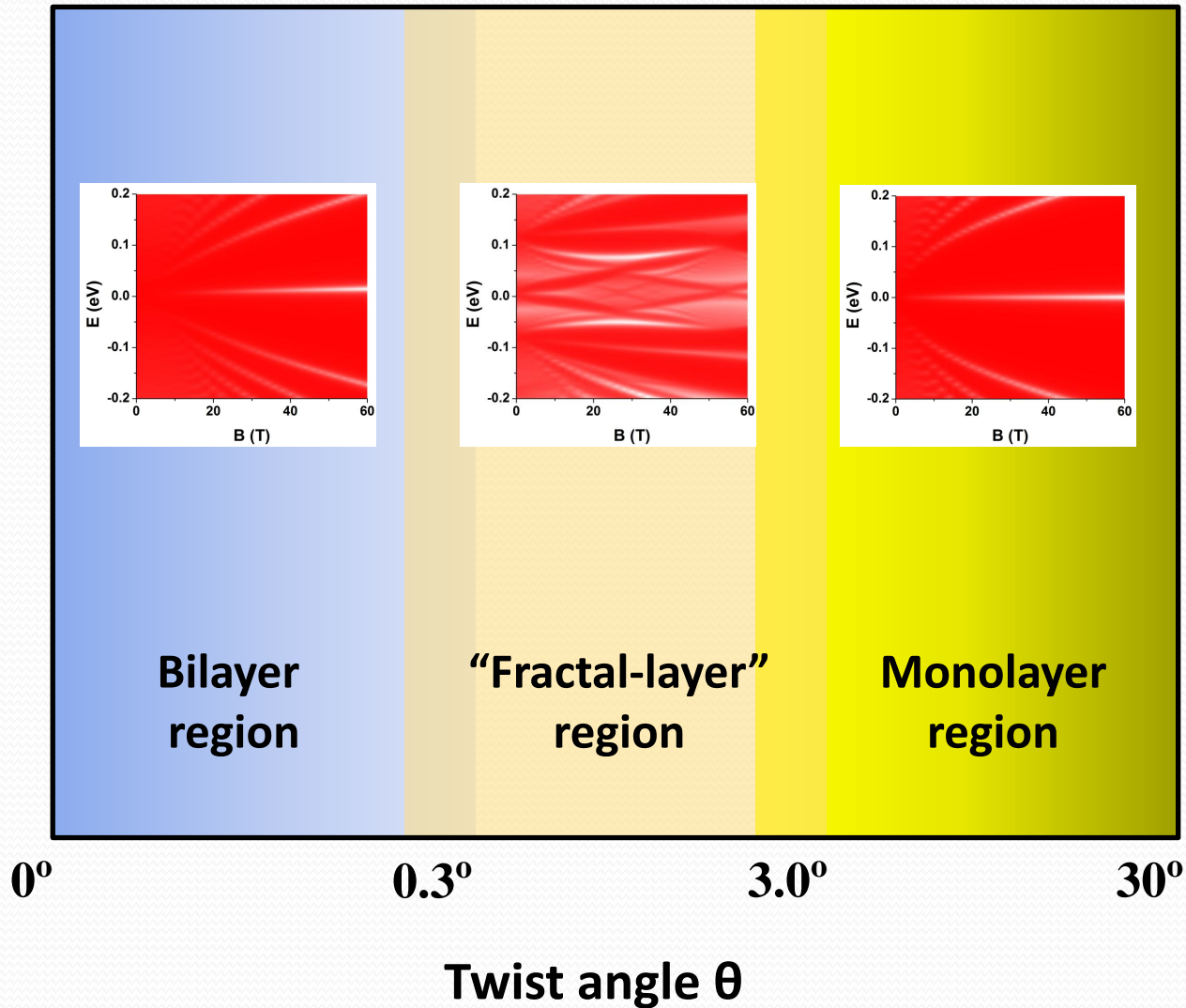




# Landau Levels as a Function of Twist Angle

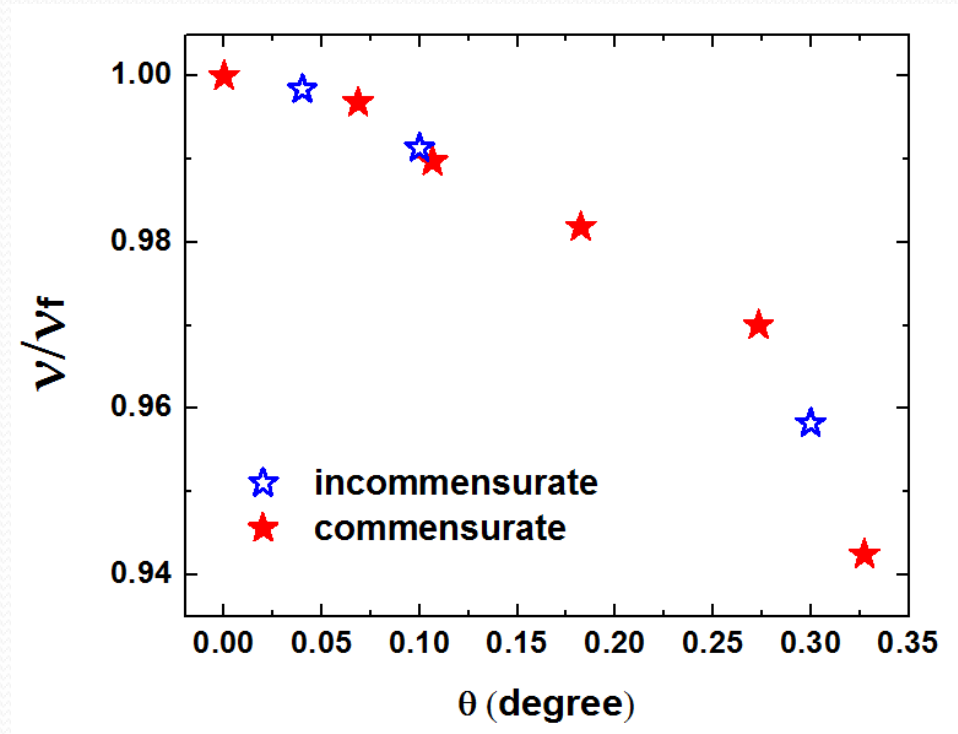
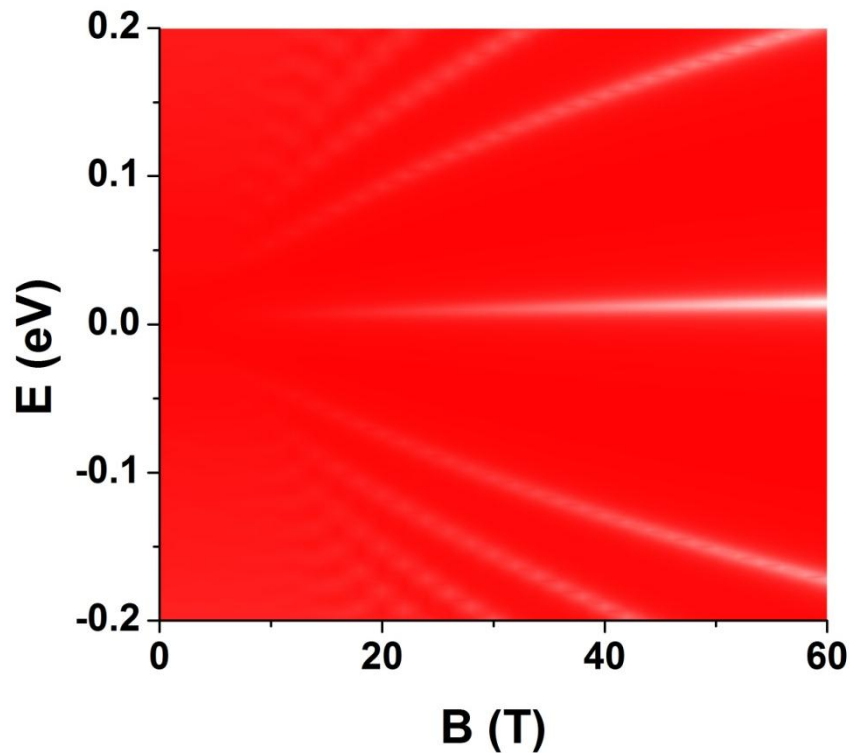


# Laudau Levels of Twisted Bilayer Graphene



# “Near-Zero” Twist Angles: Bilayer Graphene Region

Commensurate  
 $\theta = 0.06853^\circ$

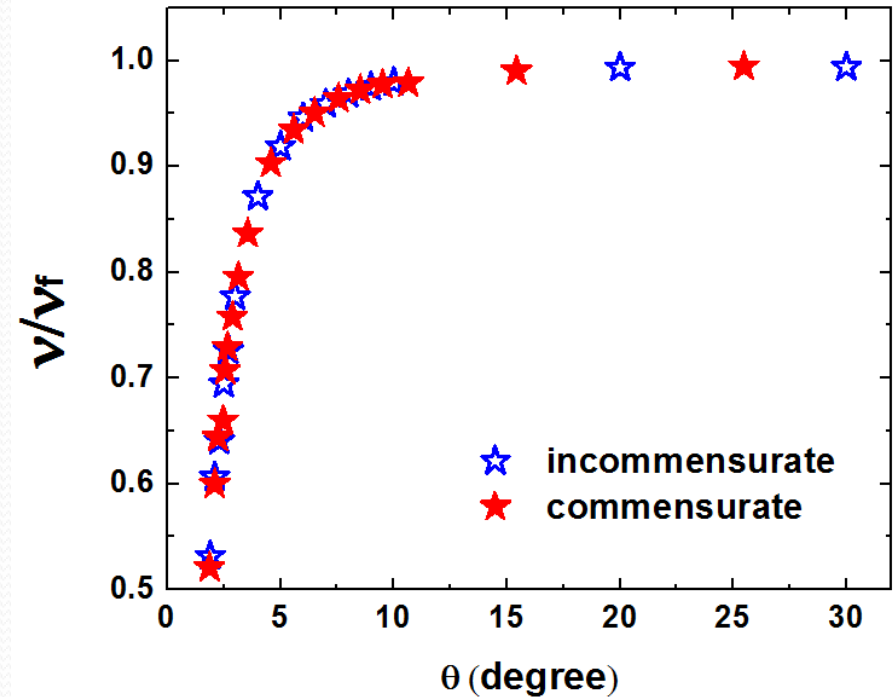
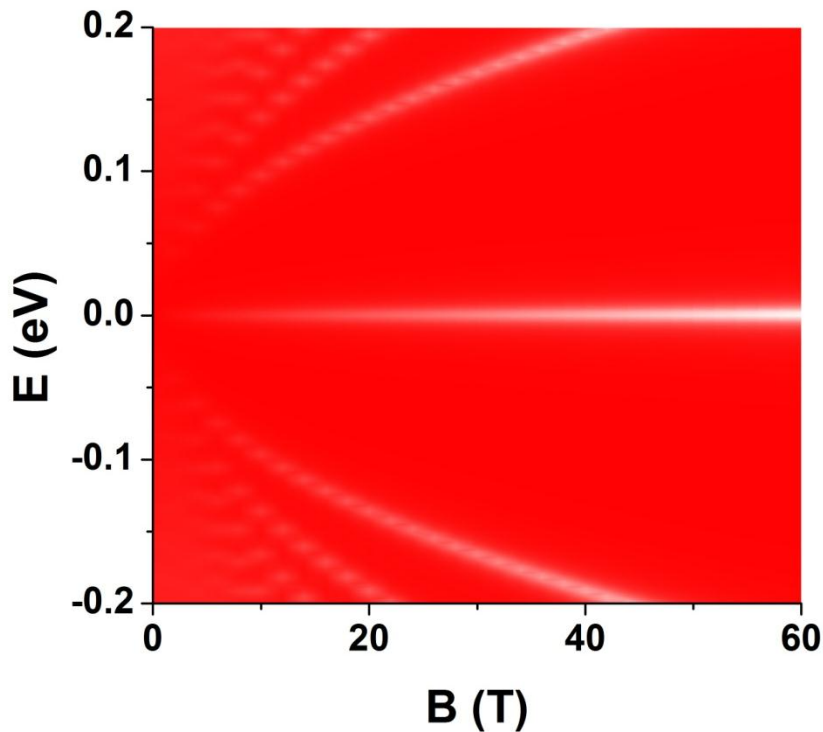


$$\Delta = \sqrt{2eBv_F^2\hbar}$$

$$E_n = \frac{\text{sgn}(n)}{\sqrt{2}} \left[ (2|n|+1)\Delta^2 + \gamma_1^2 - \sqrt{\gamma_1^4 + 2(2|n|+1)\Delta^2\gamma_1^2 + \Delta^4} \right]^{1/2}$$

# Large Twist Angles: Monolayer Graphene Region

Commensurate  
 $\theta = 7.56507^\circ$



$$E_n = \text{sgn}(n)\Delta\sqrt{|n|}$$

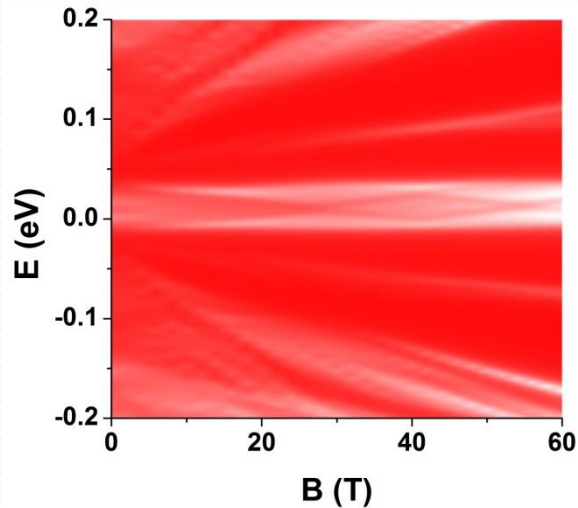
$$\Delta = \sqrt{2eBv_F^2\hbar}$$



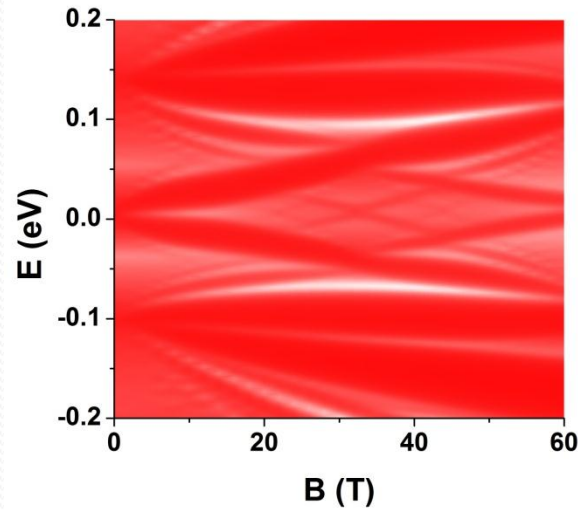
# Small Twist Angles: “Fractional-layer” Region

Commensurate:

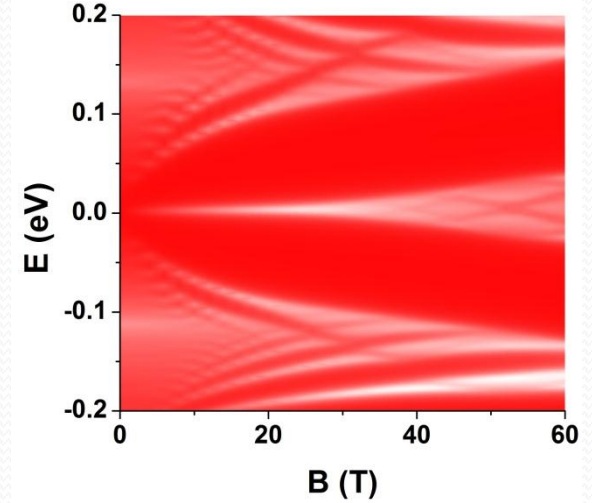
$$\theta = 1.06689^\circ$$



$$\theta = 1.64996^\circ$$

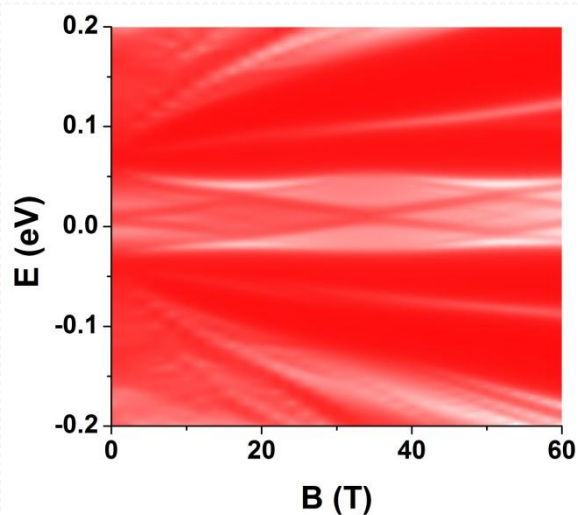


$$\theta = 2.56292^\circ$$

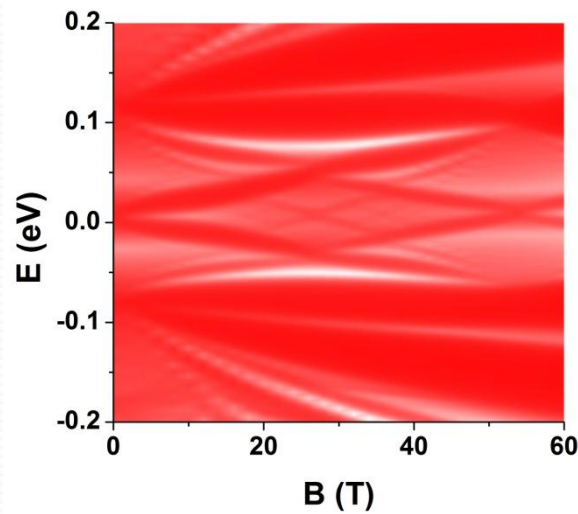


Incommensurate:

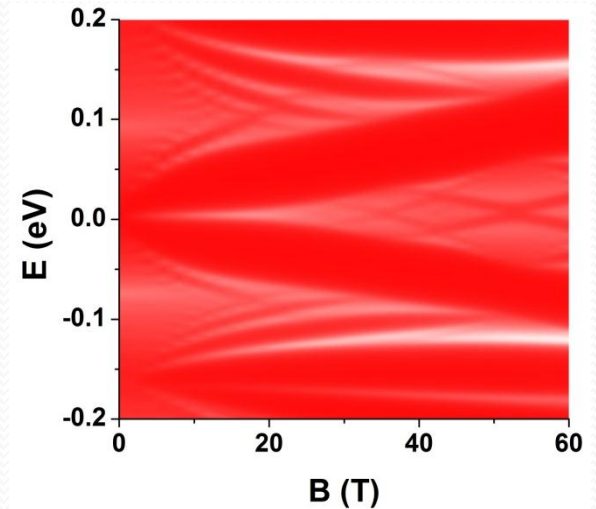
$$\theta = 1.2^\circ$$



$$\theta = 1.5^\circ$$

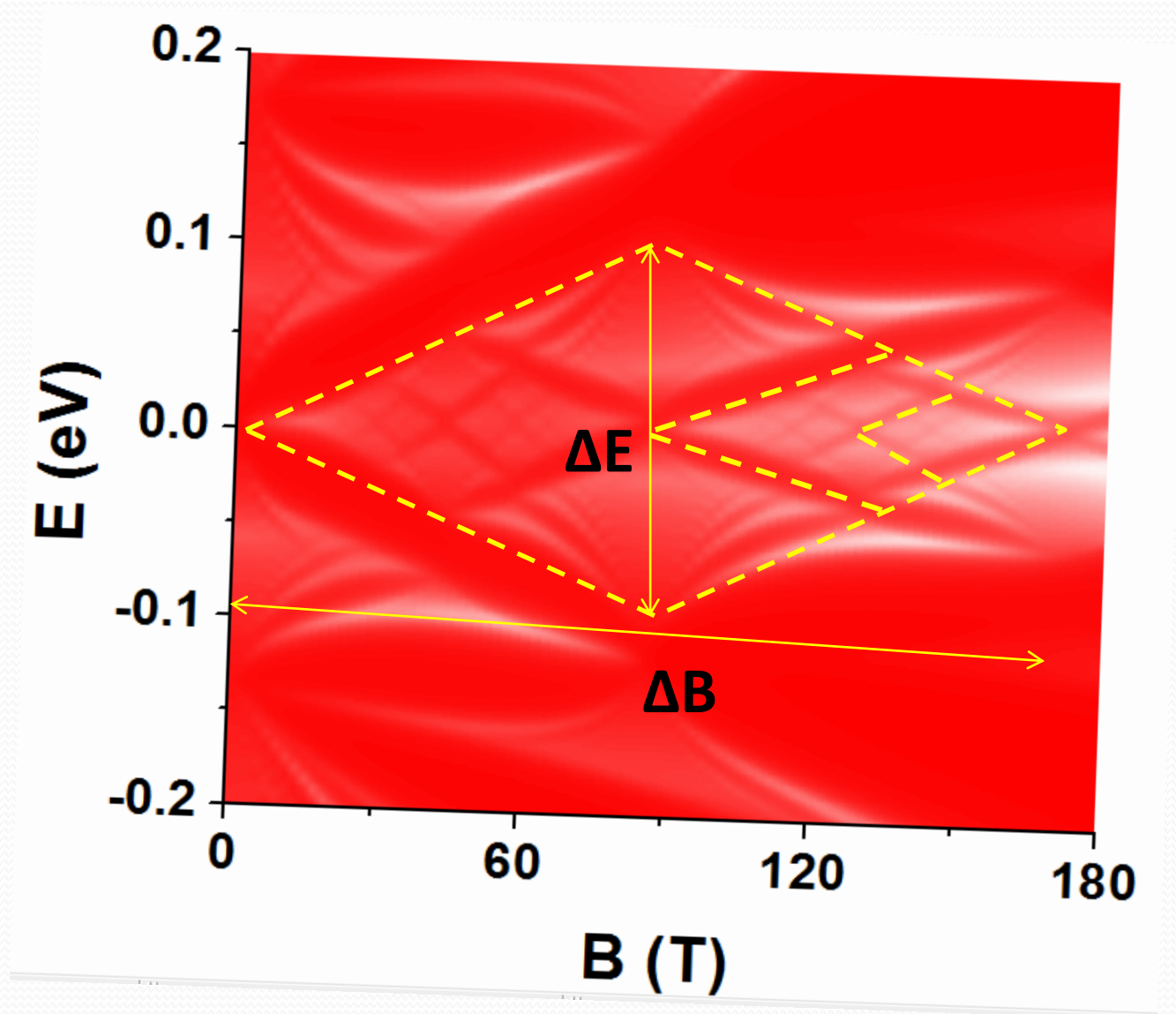


$$\theta = 2.1^\circ$$



# Fractal Spectra

Commensurate:  $\theta=1.8901^\circ$



# Conclusion

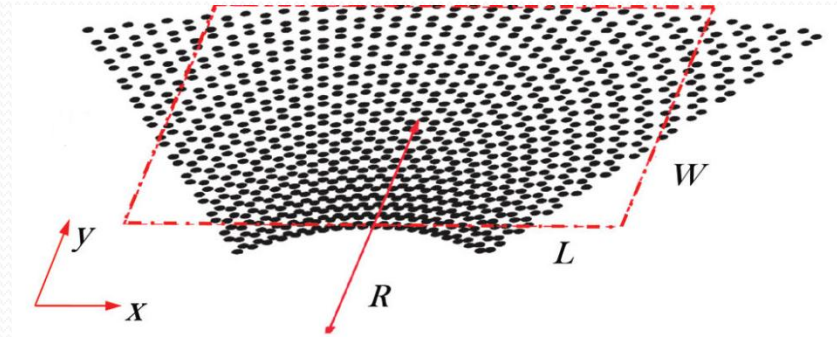
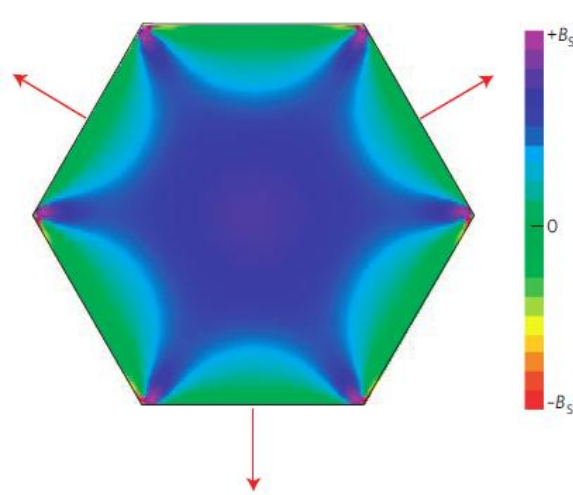
- Local LLs are independent of the lattice position for both commensurate and incommensurate twist angles in low magnetic field.
- The behavior of LLs can be classified into three regions depending on the twist angle.
- LLs in the regions of near-zero and large twist angles are characterized by a renormalized Fermi velocity as bi- and mono-layer graphene, respectively.
- In between, LLs show a complex fractional-layer behavior (Hofstadter butterfly) in a reasonably low magnetic field.



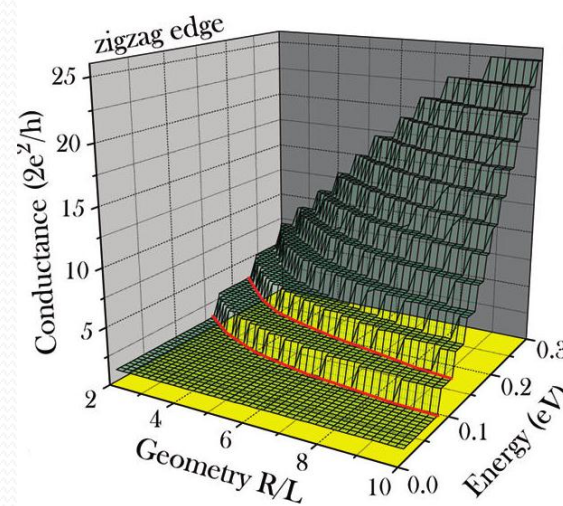
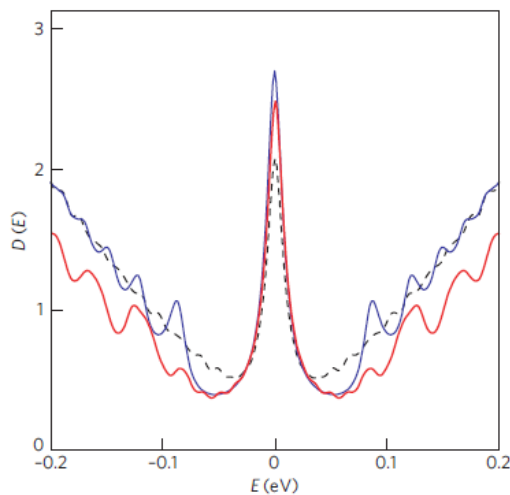
# **Pseudo Magnetic Field in Graphene Nanobubble**



# Strain-Induced Pseudo-Magnetic Fields



**Strain field**



$$\vec{A} = \frac{\beta}{a_0} \begin{pmatrix} u_{xx} - u_{yy} \\ -2u_{xy} \end{pmatrix}$$

$$\beta \approx \left. \frac{\partial \log(t)}{\partial \log(a)} \right|_{a=a_0} \approx 2-3$$

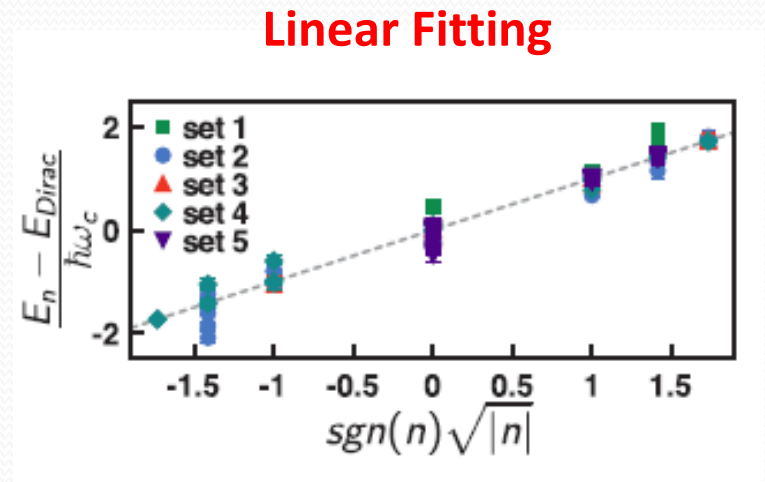
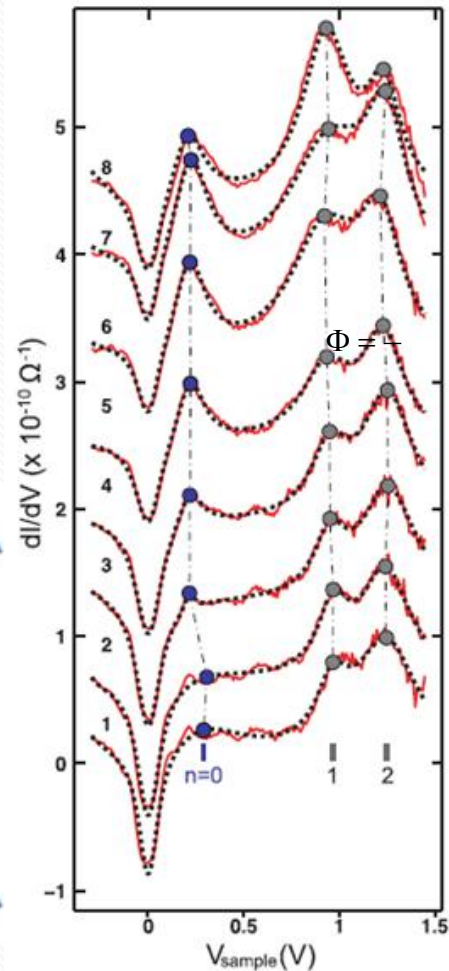
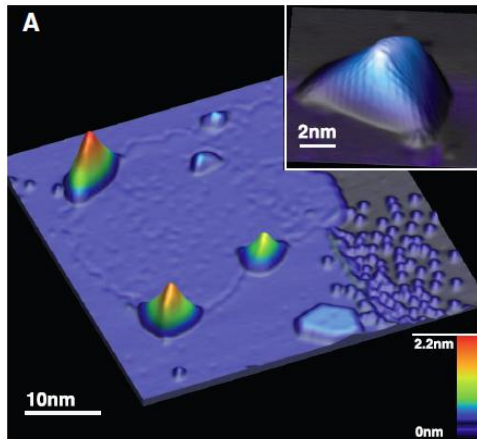
**Filling factor**

$$\nu = 2, 6, 10, \dots = 4n + 2$$

**F. Guinea, *et al.***  
**Nature Phys. 6, 30 (2010)**

**Tony Low and F. Guinea**  
**Nano Lett. 10, 3551 (2010)**

# Pseudo-Magnetic Fields in Graphene Nanobubble



**Graphene landau level**

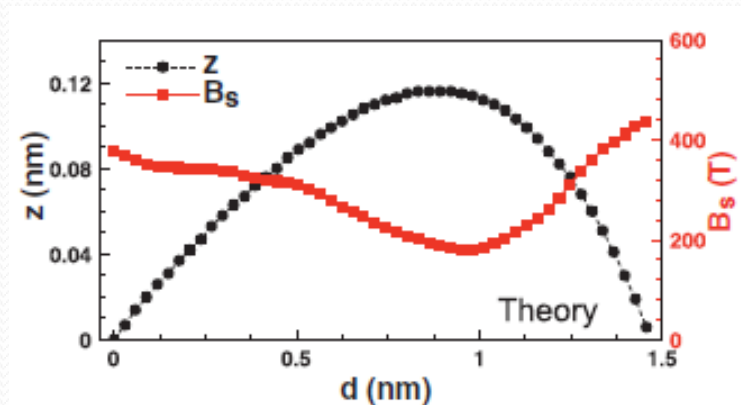
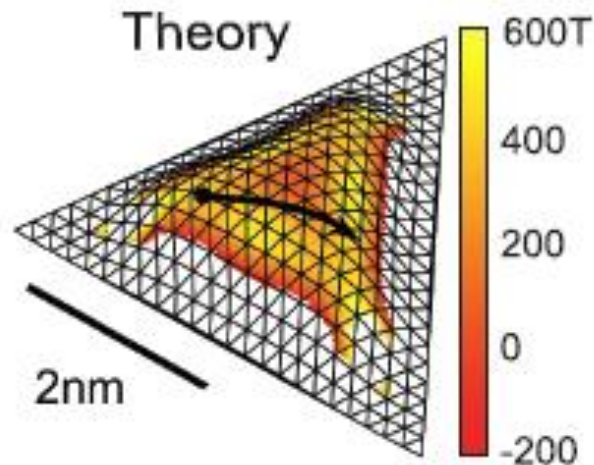
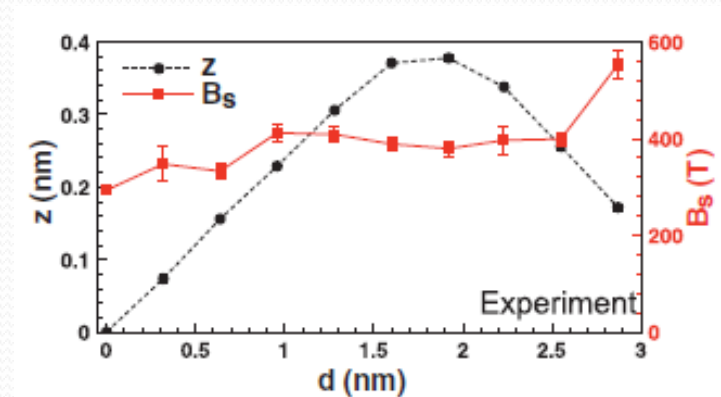
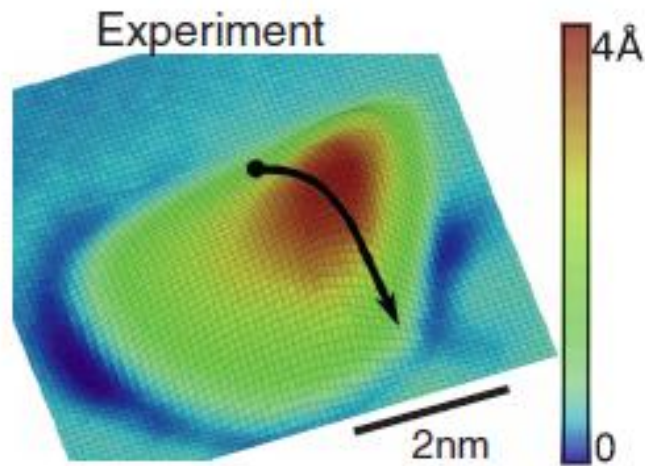
$$E_n = \text{sgn}(n)\hbar\omega_c\sqrt{|n|} + E_{\text{Dirac}},$$

$$\omega_c = \sqrt{2ehv_F^2 B_s}$$

**N. Levy, *et al.***

**Science 329, 544 (2010)**

# Strain-Induced Pseudo-Magnetic Fields ?



N. Levy, *et al.*  
Science 329, 544 (2010)

# Is This Really the “Strain-Induced” Pseudo–Magnetic Field Effect ?

➤ Strain effect ?

➤ Curvature effect ?

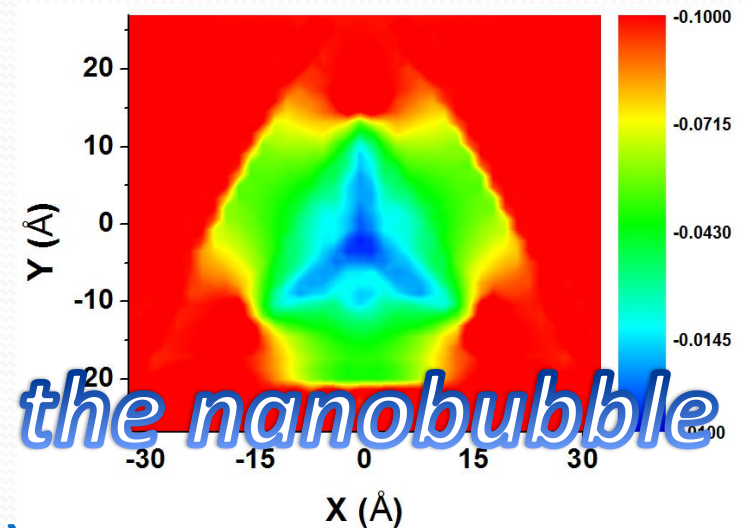


# Strain Map of 2D Graphene Nanobubble

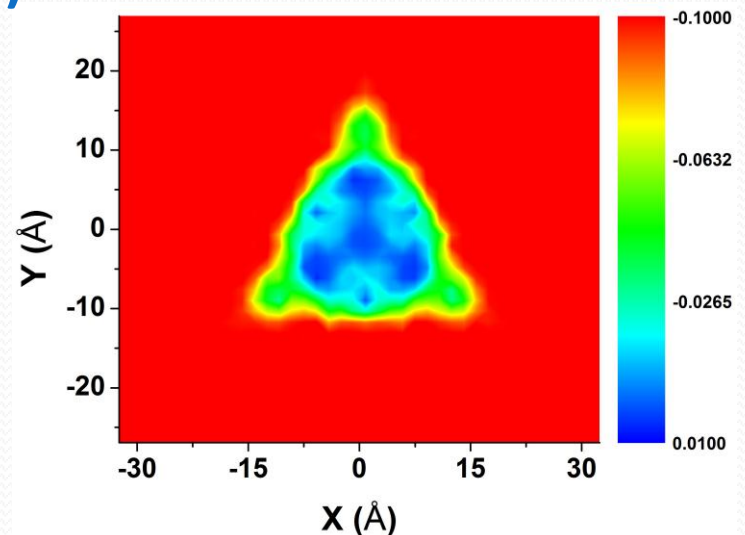
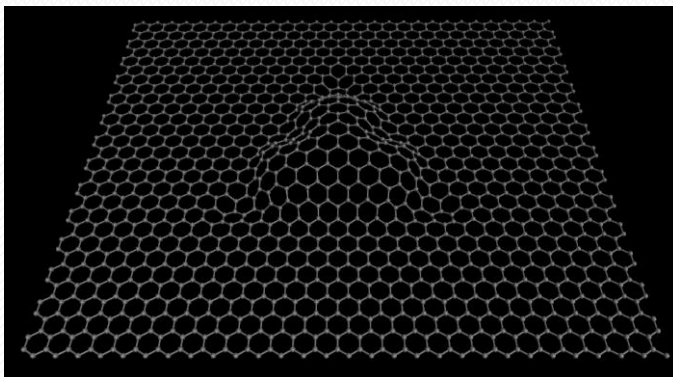
## ➤ Case one: fix z (10% strain)



*Almost strain free in the nanobubble*



## ➤ Case two: relax all (10% strain)



# 1D Graphene Nanobubble

